

The Nuclear Symmetry Energy in Astrophysics

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5. *The super-nova process*

We have tentatively suggested that the super-nova process represents the transition of an ordinary star into a neutron star. If neutrons are produced on the surface of an ordinary star they will "rain" down towards the center if we assume that the light pressure on neutrons is practically zero. This view explains the speed of the star's transformation into a neutron star. We are fully aware that our suggestion carries with it grave implications regarding the ordinary views about the constitution of stars and therefore will require further careful studies.

W. BAADE

F. ZWICKY

Mt. Wilson Observatory and

California Institute of Technology, Pasadena.

May 28, 1934.

neutron star basics



$$R \sim 1 \text{ fm} \cdot (10^{57})^{1/3} \sim 10 \text{ km}$$

$$\text{redshift} \equiv \frac{\Delta\lambda}{\lambda} \approx 0.3$$

$$\frac{GMm_{\text{H}}}{R} \approx 200 \text{ MeV}$$

NEUTRON STAR MODELS

A. G. W. CAMERON

Atomic Energy of Canada Limited, Chalk River, Ontario, Canada

Received June 17, 1959

ABSTRACT

Previous models of neutron stars were constructed with the assumption that the equation of state of a neutron gas is that of non-interacting Fermi gas. Such models have a maximum observable mass of about 0.7 solar mass. In fact, the potential energy of a neutron gas depends on the density; this introduces additional terms into the equation of state. A revised equation of state has been derived which makes use of a mean nuclear potential recently given by T. H. R. Skyrme. Twenty neutron star models have been constructed by integrating the general relativistic equations of hydrostatic equilibrium of the neutron gas. The results show that there is an upper limit to the observable mass of about 2 solar masses; the corresponding upper limit to the proper mass is about 3 solar masses. There is a lower limit to each of these masses of about 0.05 solar mass, below which the neutron star is unstable against transformation into an iron star. The radii of these neutron stars lie in the range 7-9 km. A qualitative discussion of the effects of transformation of neutrons into hyperons at very high densities is given.

$$E = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2 + 5kx$$

G

BERNOULLI

$$\frac{p}{\rho} + gz + \frac{v^2}{2} = H$$

GLIDER RIDING A THERMAL

Shock Wave

Shock Wave

VELOCITY PRESSURE

LOW HIGH

VELOCITY PRESSURE

LOW HIGH

VENTURI

GRAVITY

Introduction

From the microscopic to the macroscopic

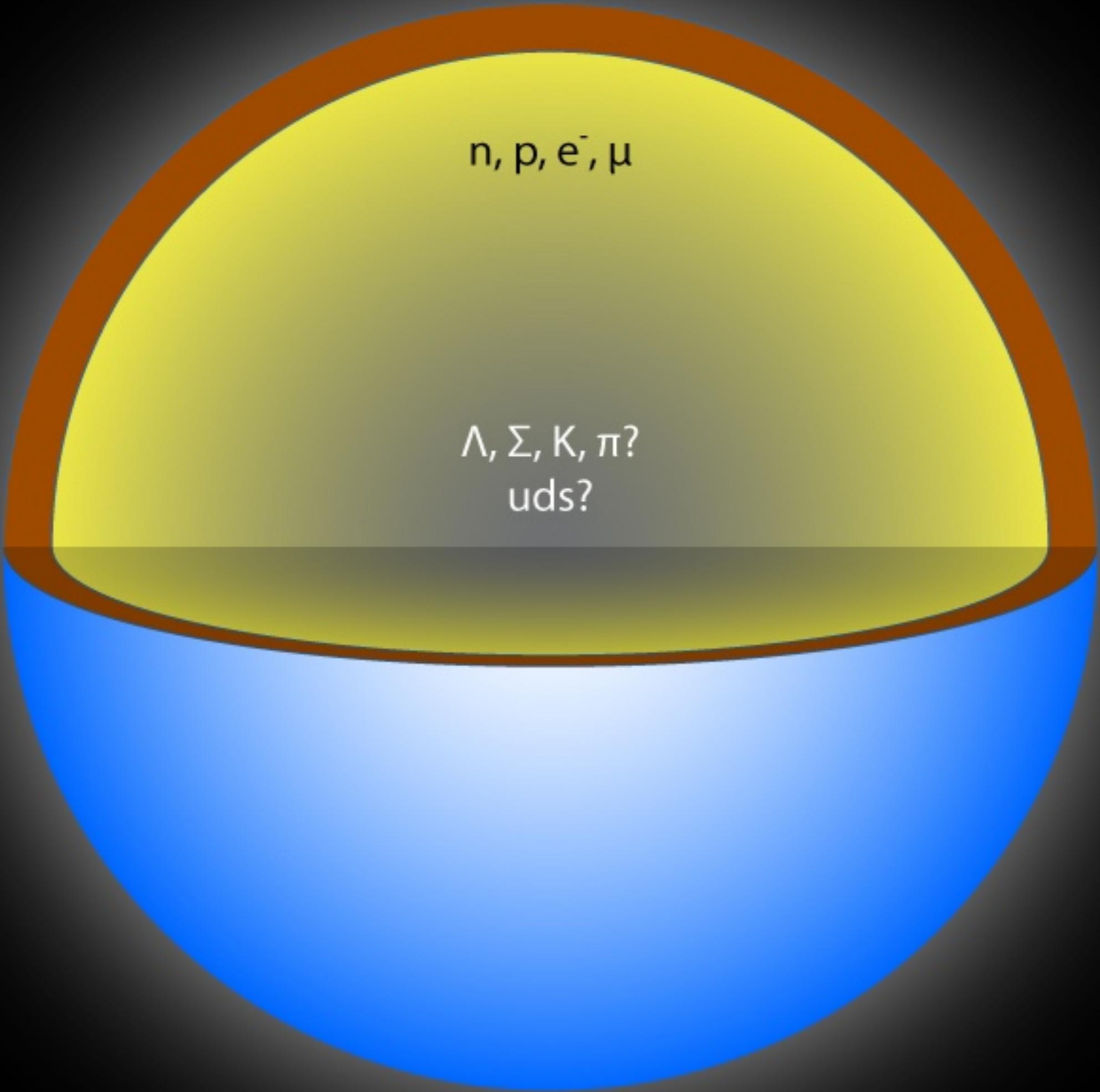
Constraints from astronomy

Pulsar timing

X-ray burst spectroscopy

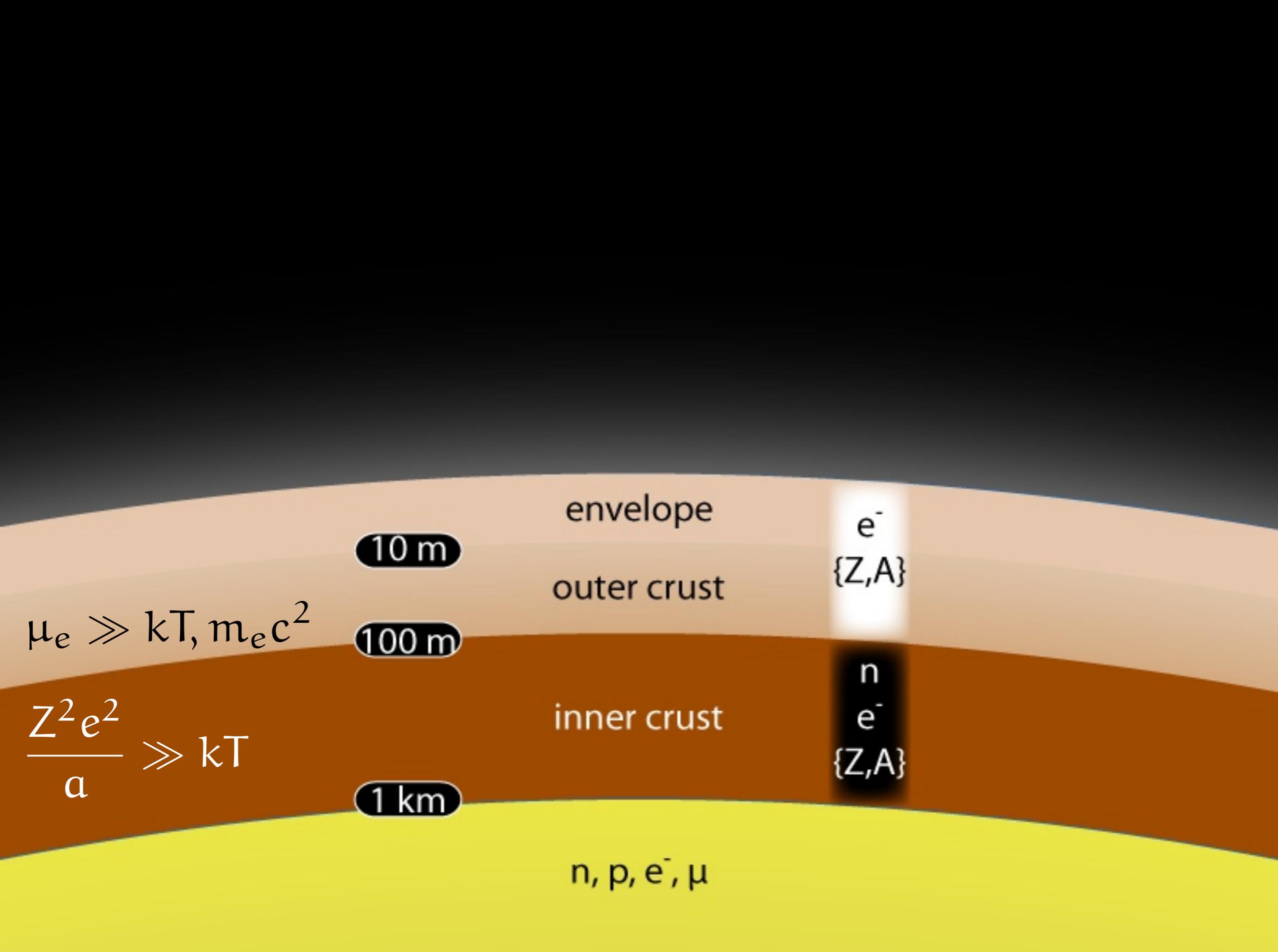
Neutron star cooling

Conclusion



n, p, e^-, μ

$\Lambda, \Sigma, K, \pi?$
 $uds?$



envelope

10 m

e^-
 $\{Z,A\}$

outer crust

100 m

$$\mu_e \gg kT, m_e c^2$$

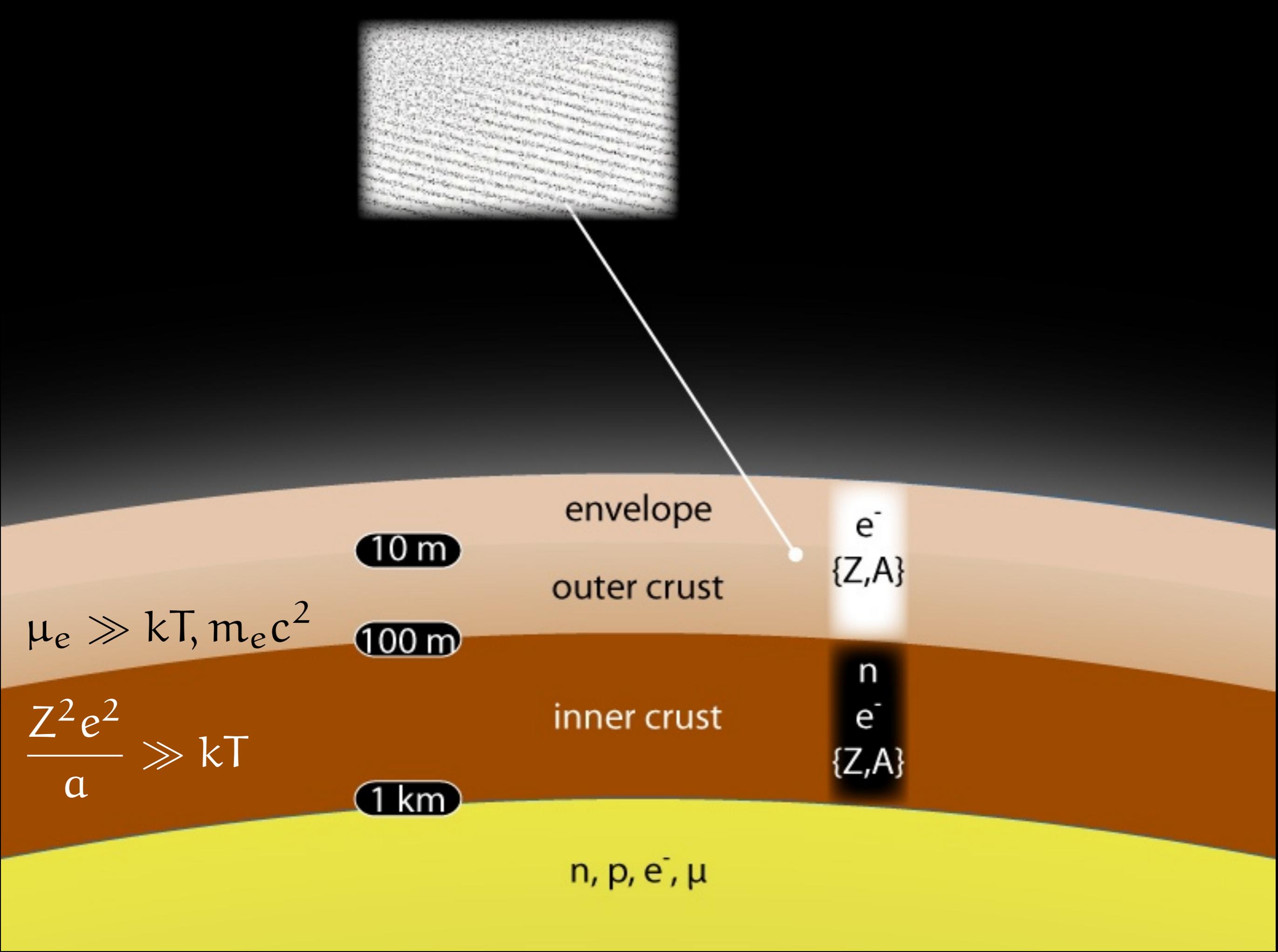
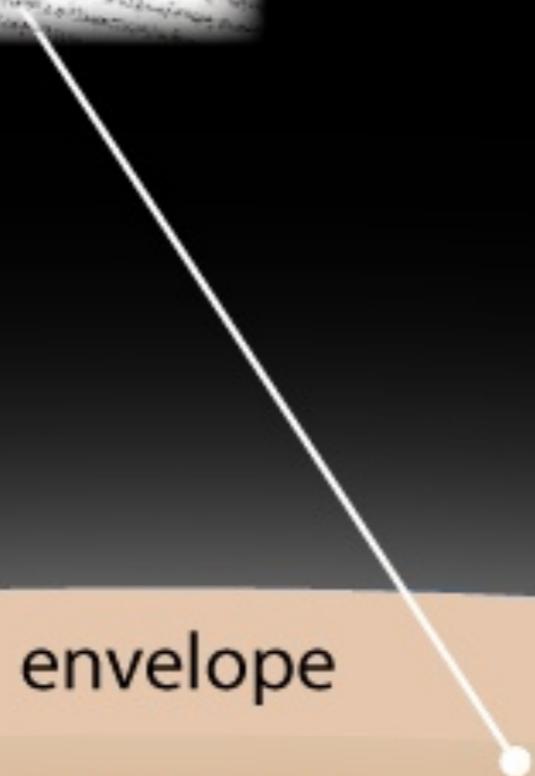
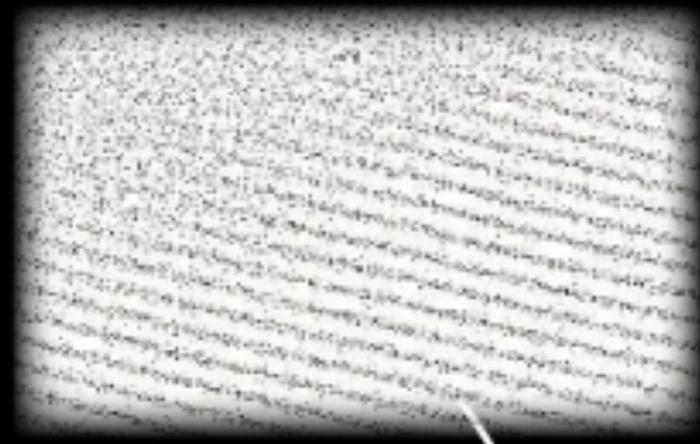
inner crust

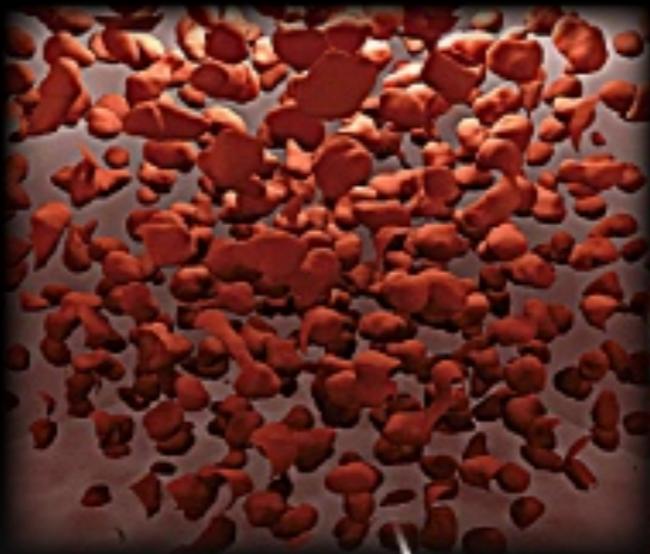
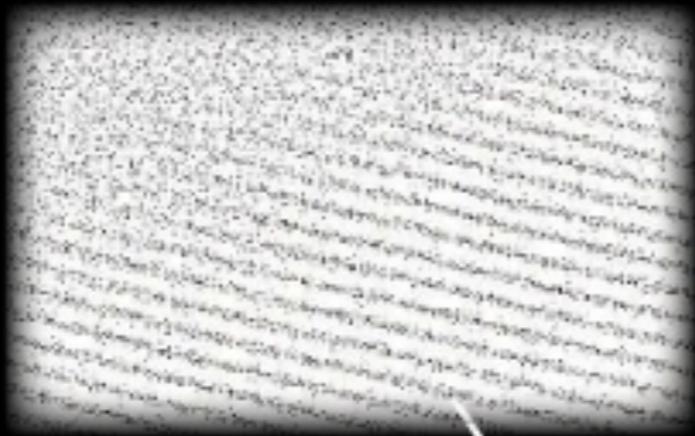
n
 e^-
 $\{Z,A\}$

$$\frac{Z^2 e^2}{a} \gg kT$$

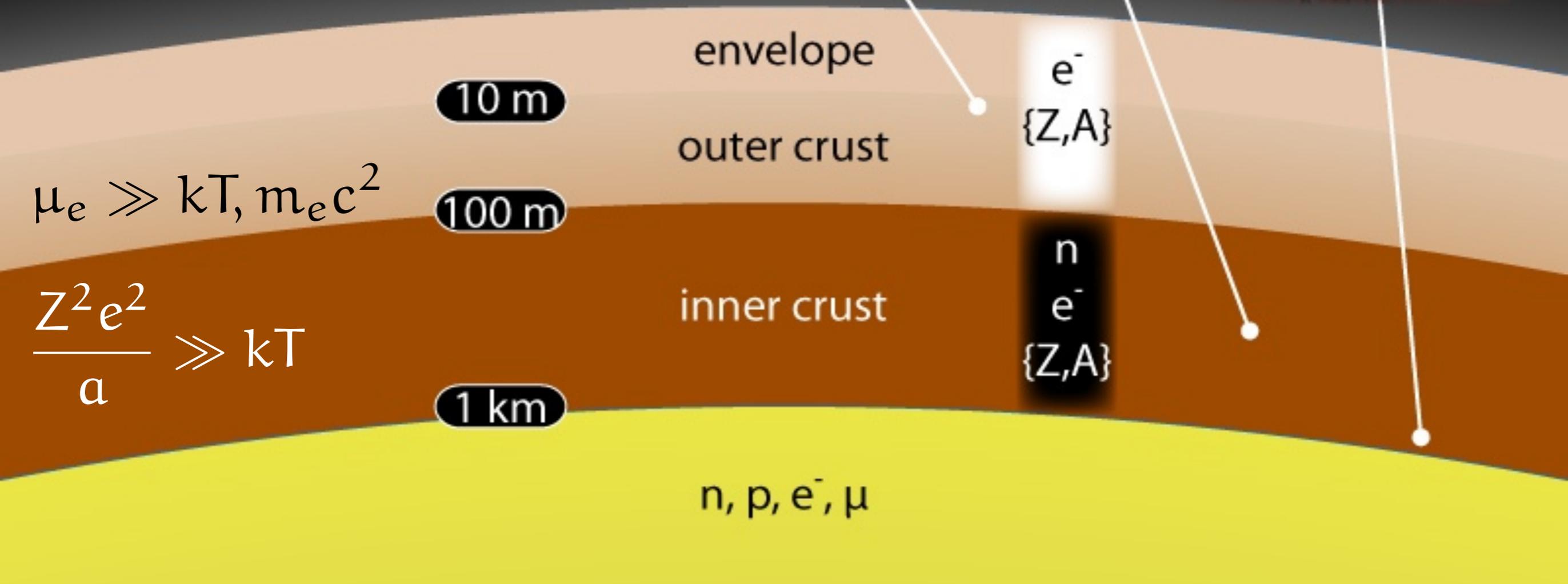
1 km

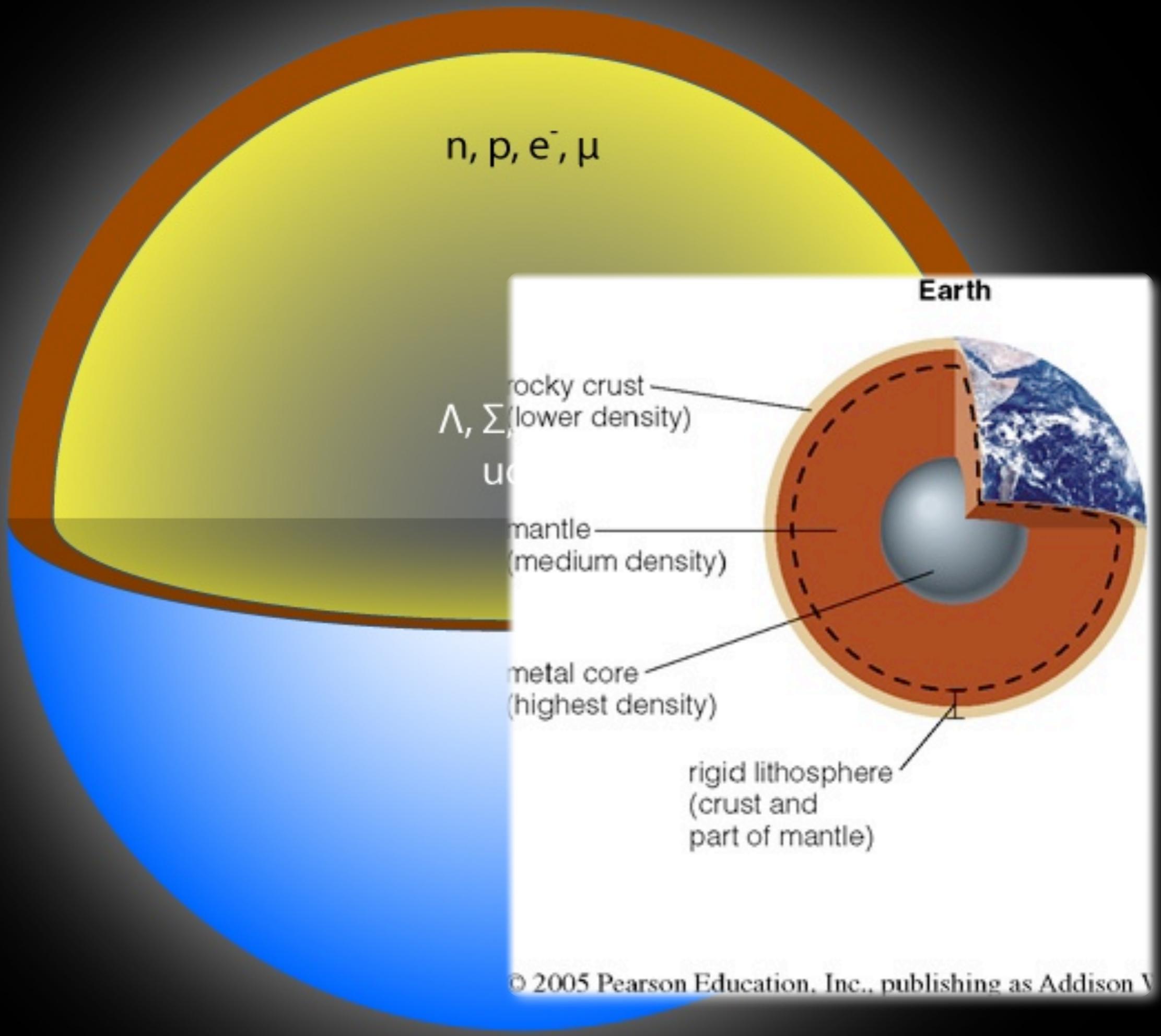
n, p, e^-, μ





Talks by Newton,
Gandolfi





n, p, e^-, μ

Earth

rocky crust
(lower density)
 Λ, Σ
u

mantle
(medium density)

metal core
(highest density)

rigid lithosphere
(crust and
part of mantle)

Thermodynamics near saturation density

see review by Lattimer & Prakash

Given n_{pe} with $x = n_p / (n_n + n_p)$, with energy per nucleon

$$\varepsilon(n, x) = \varepsilon_S(n) + \varepsilon_A(n)(1 - 2x)^2 :$$

- in β -equilibrium,

$$\mu_e = \mu_n - \mu_p = -\frac{\partial \varepsilon}{\partial x} = 4\varepsilon_A(1 - 2x);$$

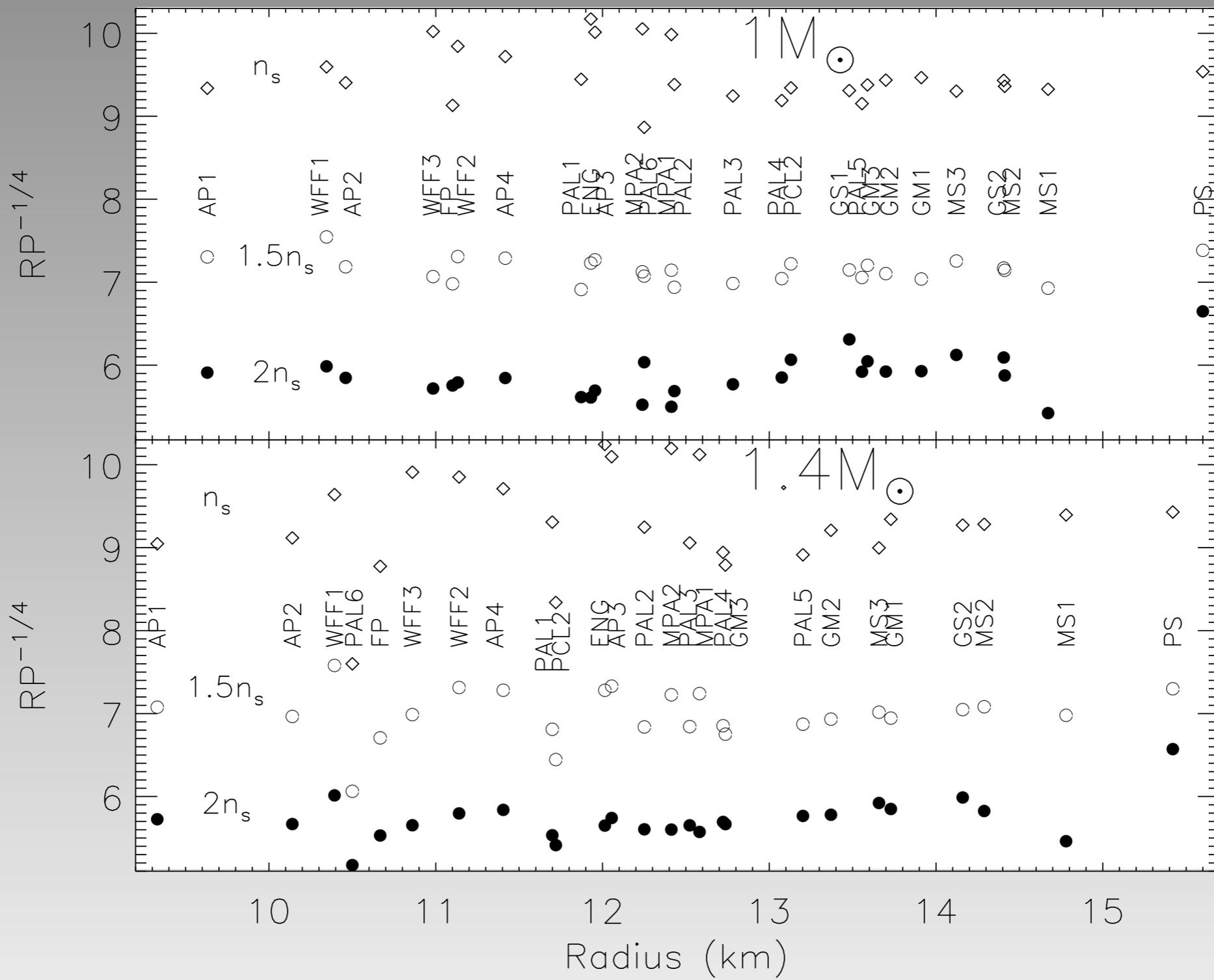
- the total pressure is

$$P = n^2 \frac{\partial \varepsilon}{\partial n} + \frac{n_e \mu_e}{4} = n(1 - 2x) \left[n \frac{\partial \varepsilon_A}{\partial n} (1 - 2x) + x \varepsilon_A \right];$$

- and upon solving $(3\pi^2 n_e)^{1/3} \hbar c = \mu_e$, the proton fraction is

$$x \approx \left[6 + \frac{3\pi^2}{64} \left(\frac{\hbar c}{\varepsilon_A} \right)^3 n \right]^{-1} \approx 0.04$$

for $\varepsilon_A = 30$ MeV.



radius correlated with pressure near saturation density

Lattimer & Prakash 2001

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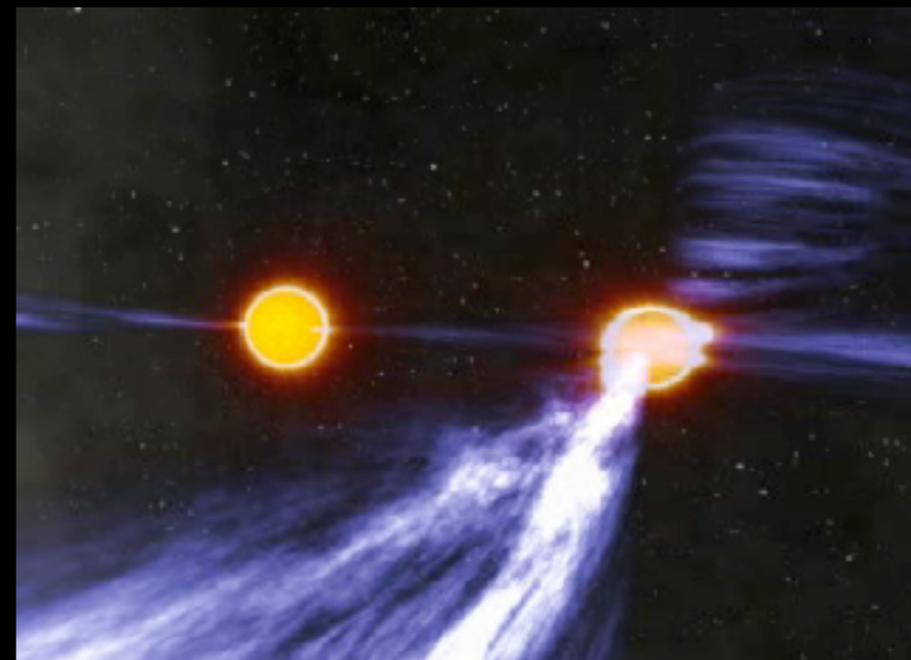
Detection: Isolated neutron stars

about 1700 pulsars detected; about 50 are in binary systems with some mass information

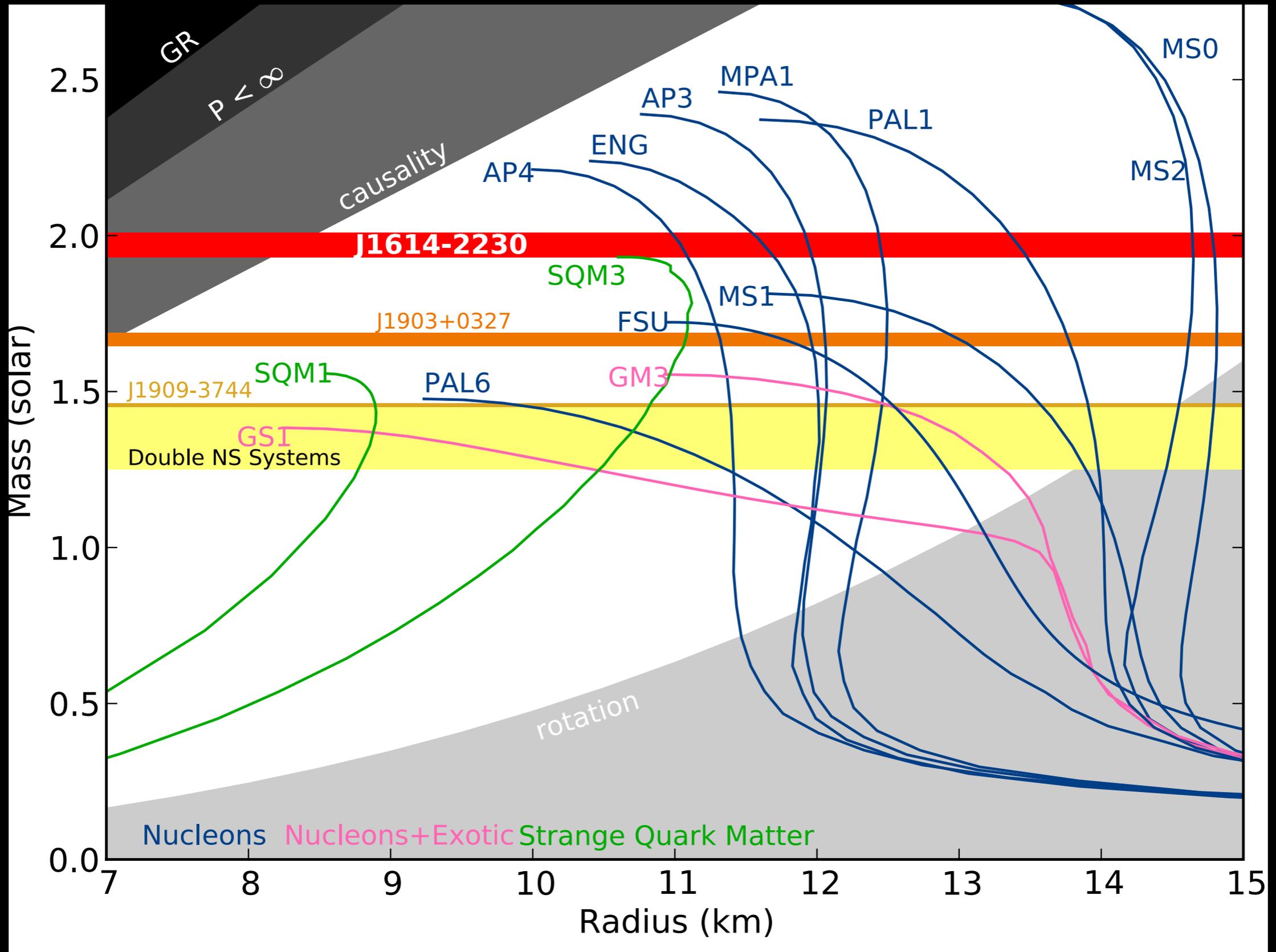
very precise mass information;

but, no radius information

PSR J0737-3039A/B

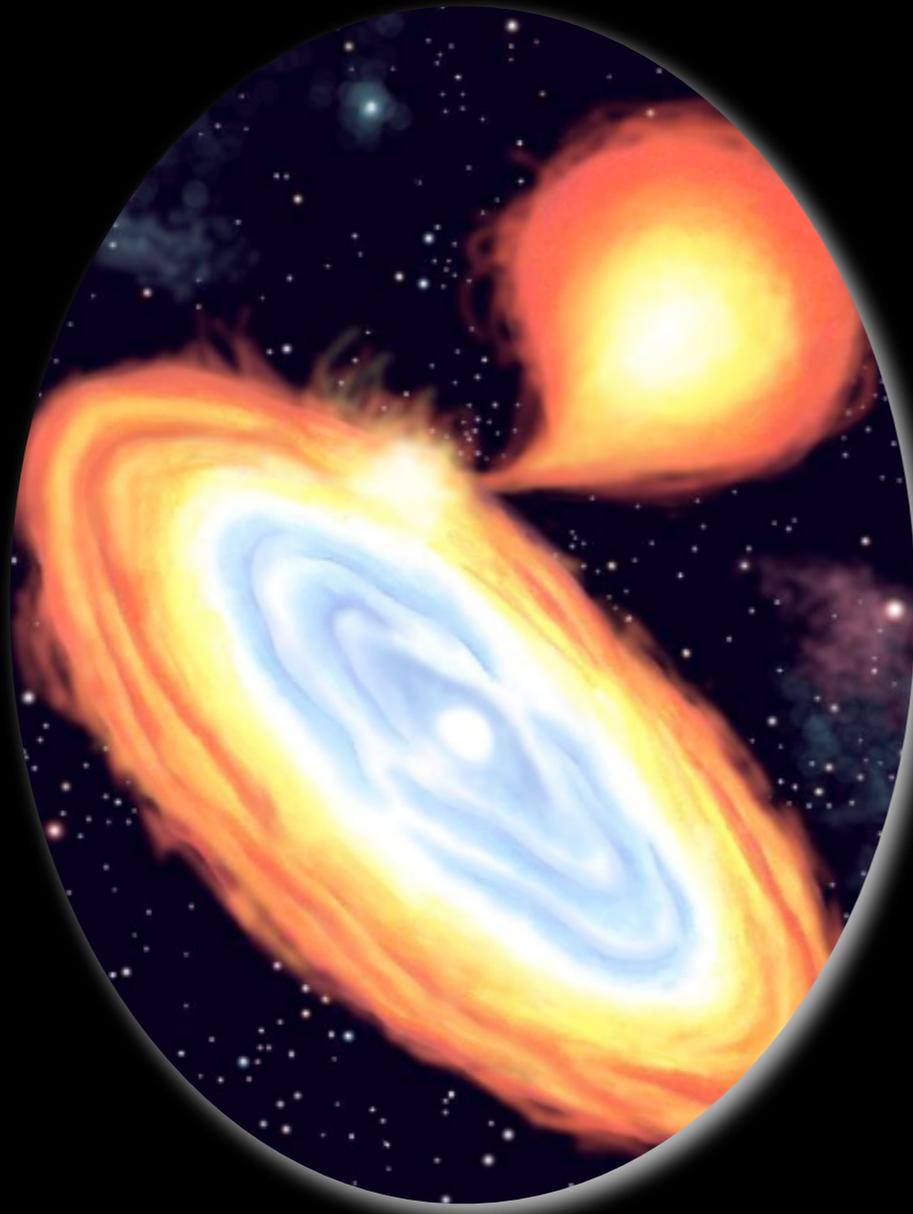


John Rowe Animation/Australia Telescope
National Facility, CSIRO



Pulsar Masses, Demorest et al. 2010

Detection: Accreting neutron stars



About 100 low-mass X-ray binaries known

rich phenomenology:

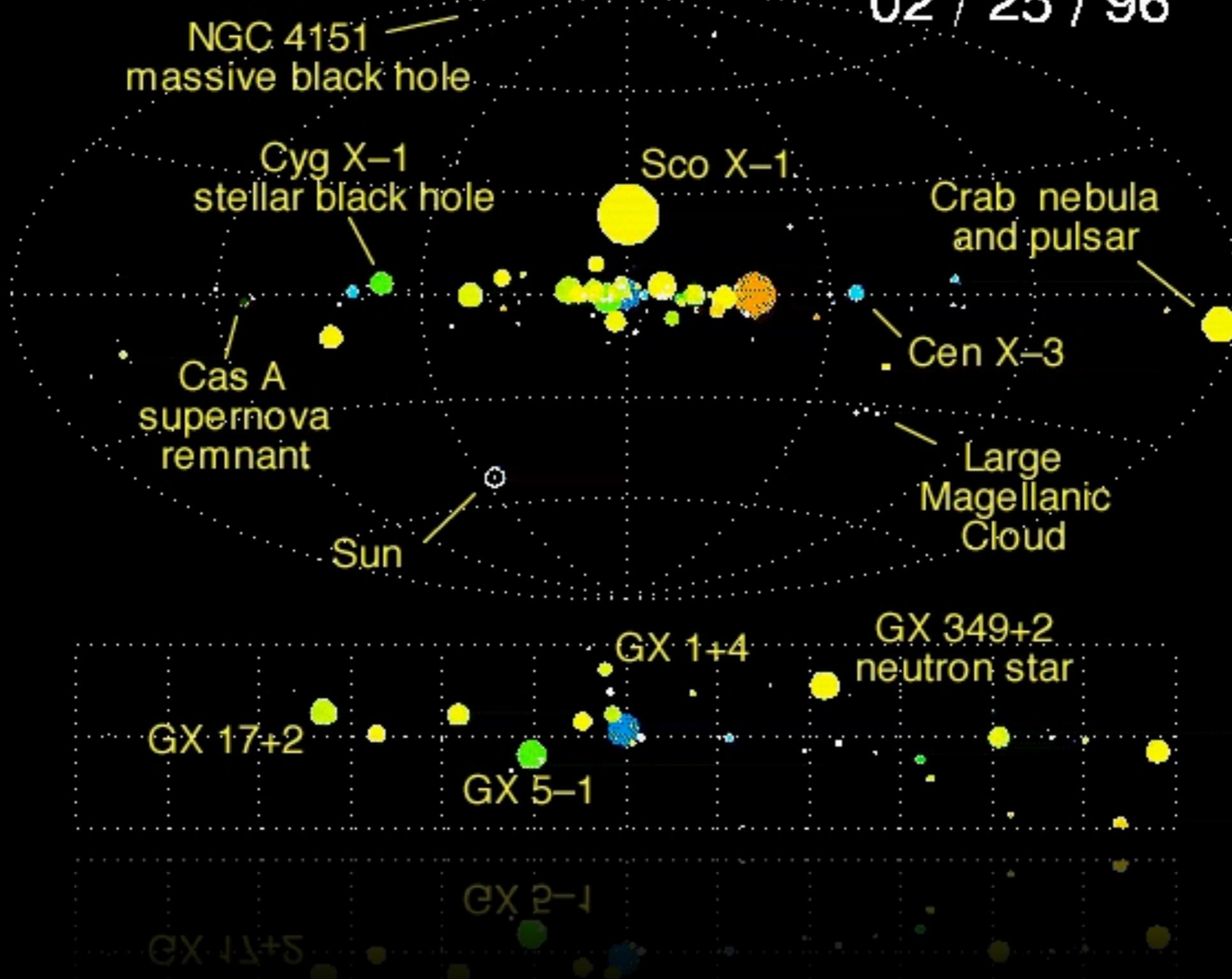
bursts, superbursts, oscillations

crust not in thermal equilibrium with core; but

measurements not as precise

Some X-ray Landmarks

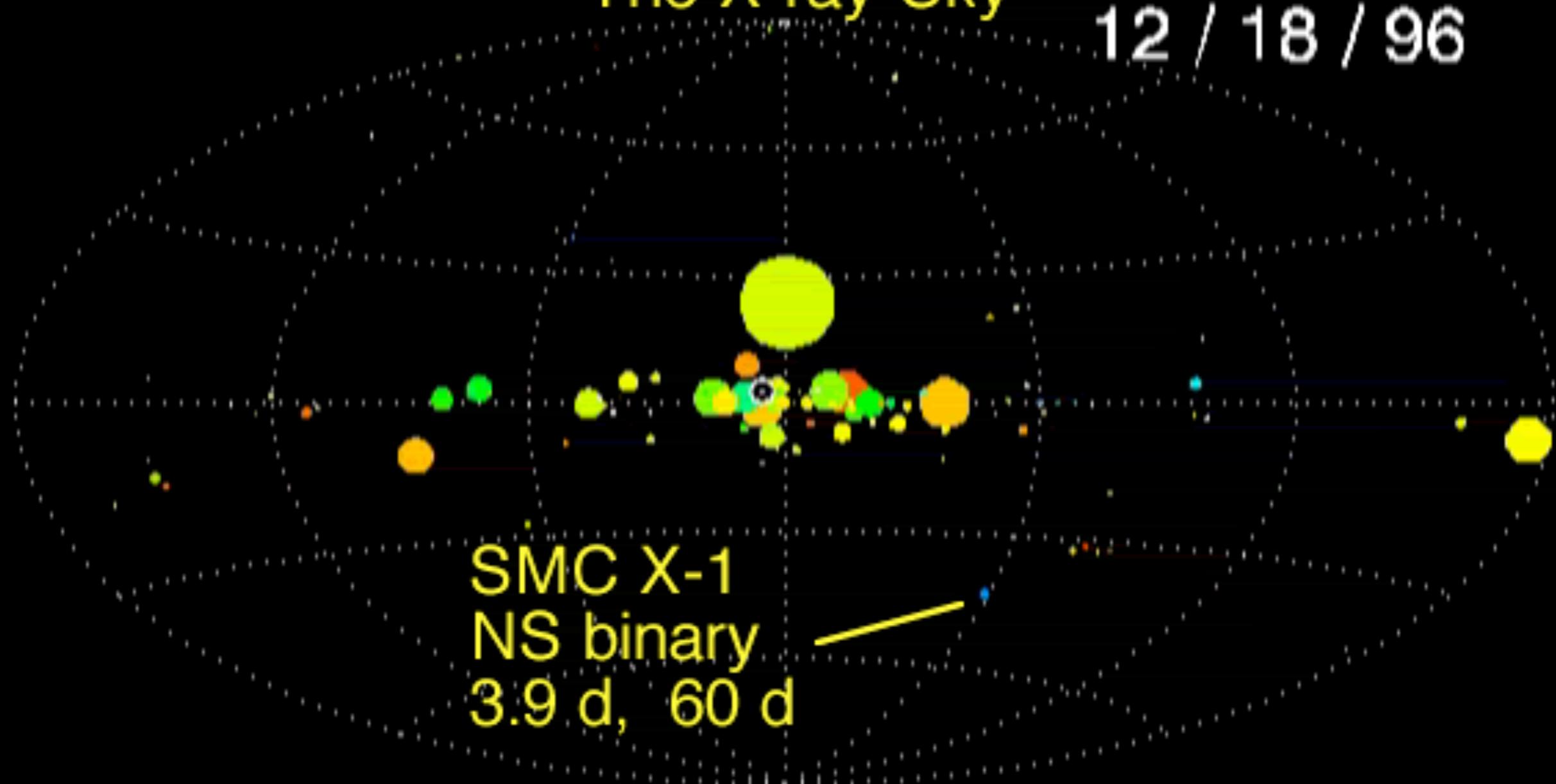
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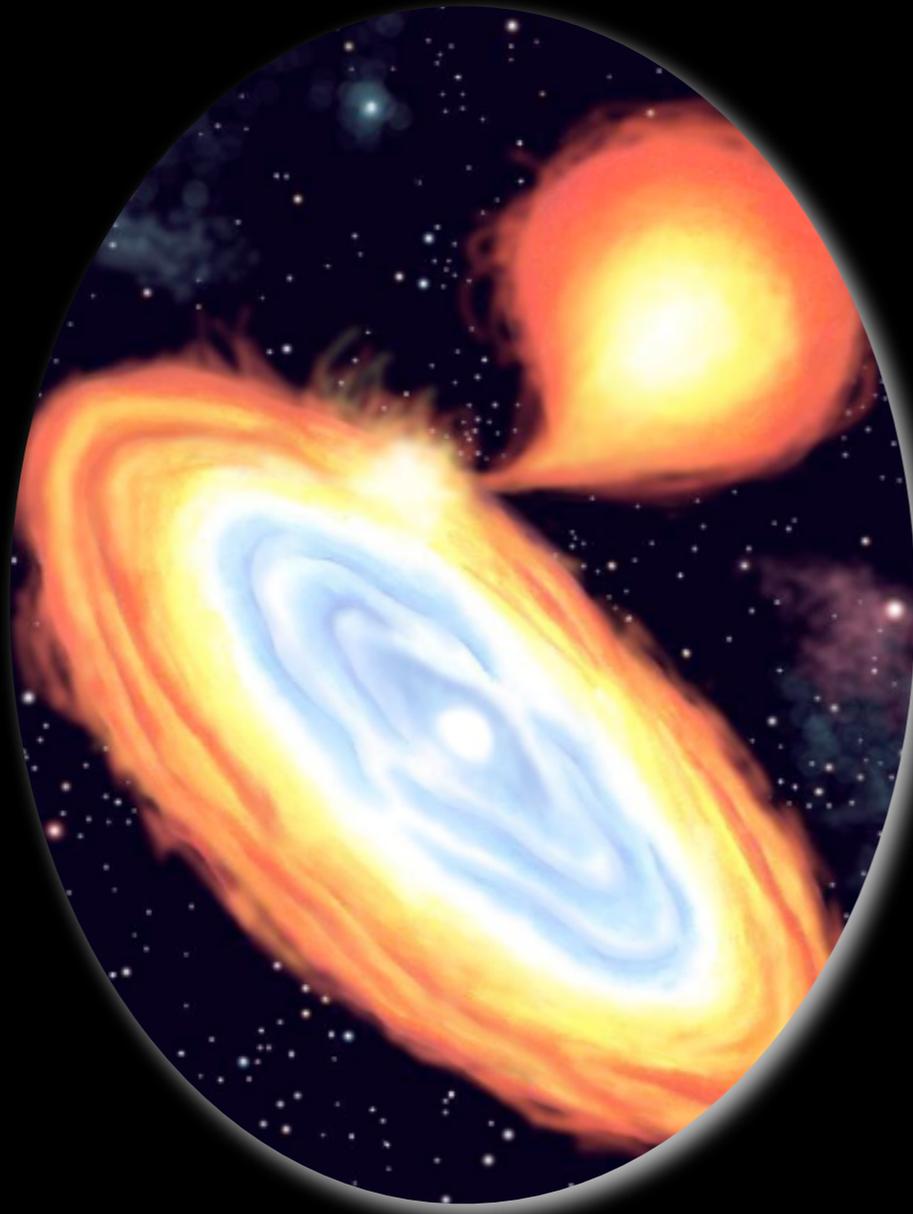
The X-ray sky, courtesy RXTE/ASM monitor team

The X-ray Sky

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Detection: Accreting neutron stars



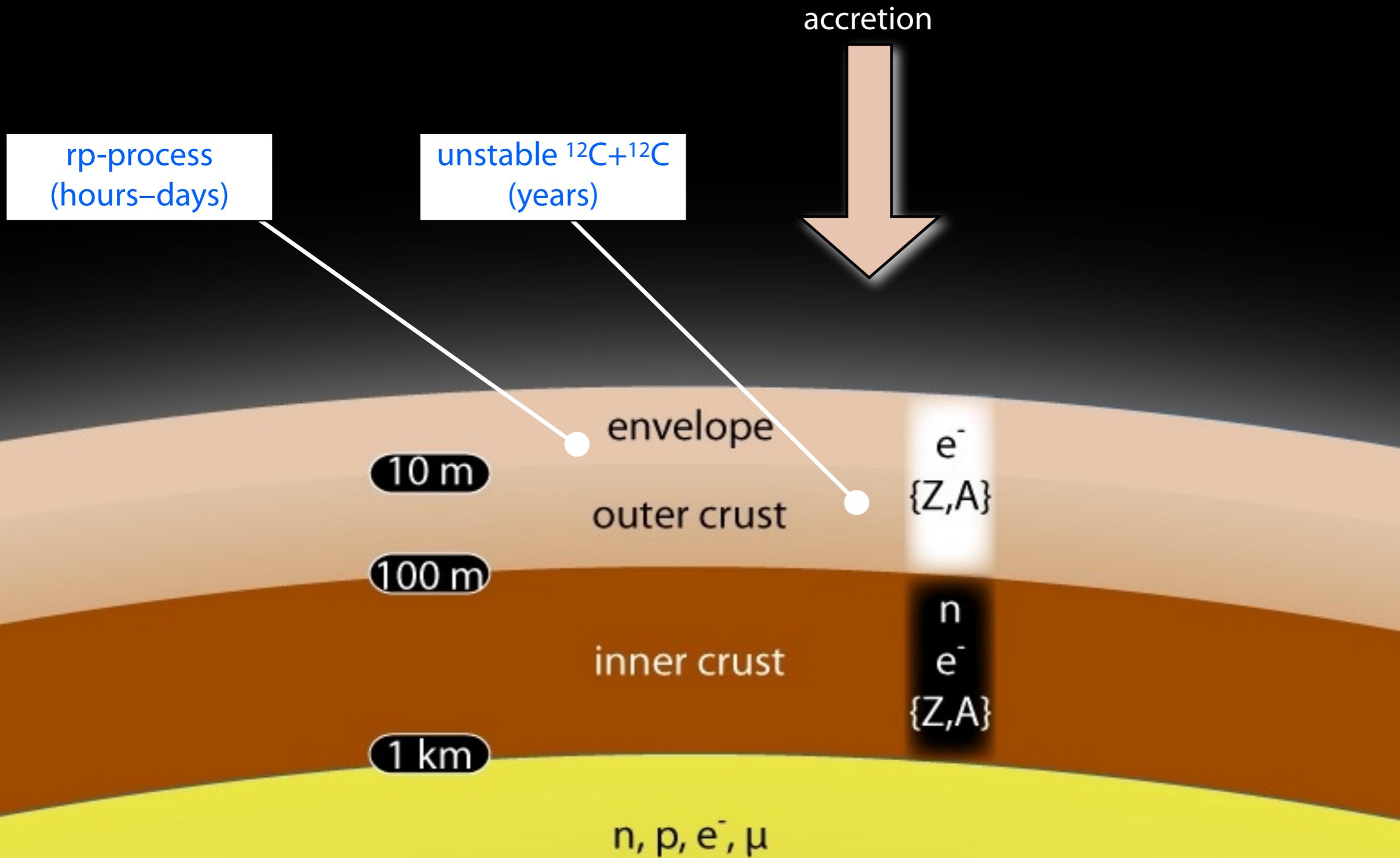
About 100 low-mass X-ray binaries known

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bursts, superbursts,
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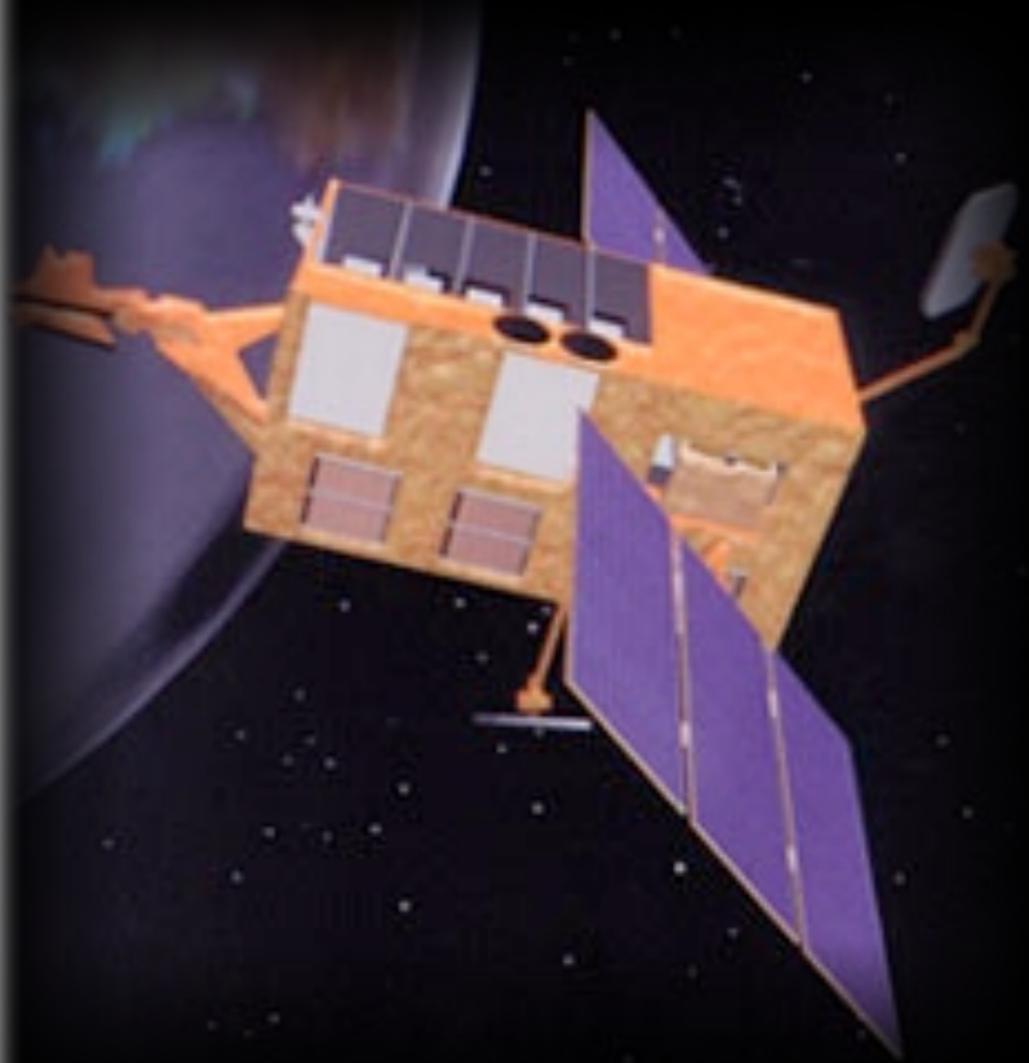
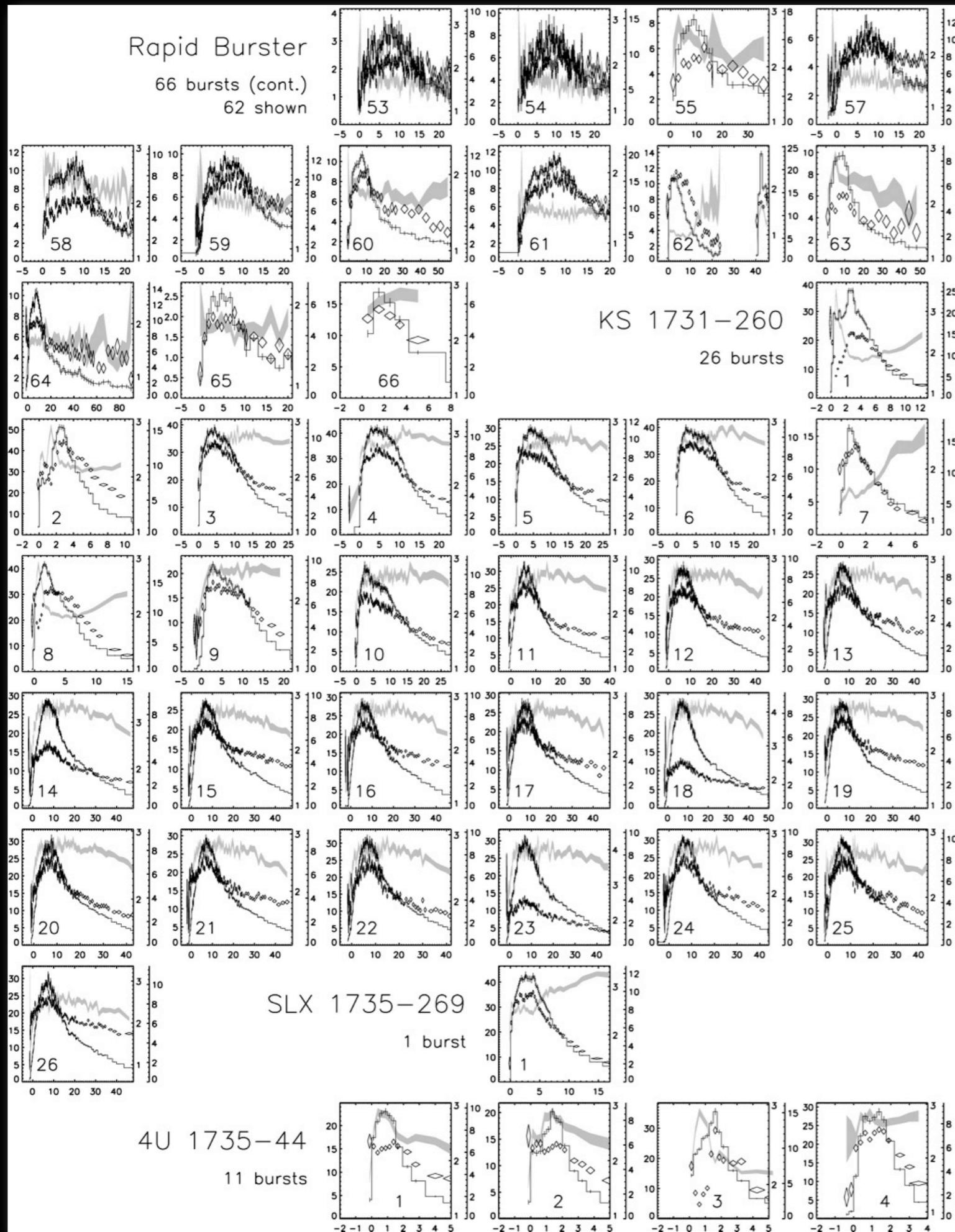
crust not in thermal equilibrium
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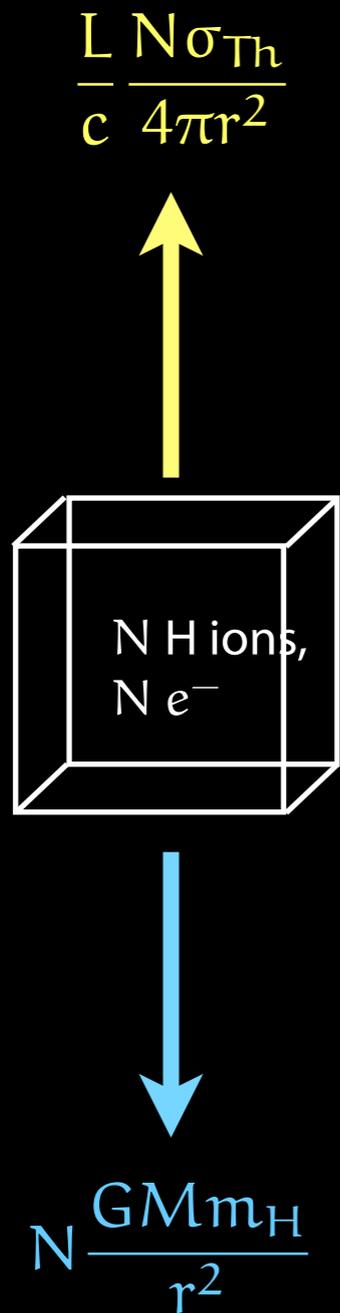


Galloway et al. 2008

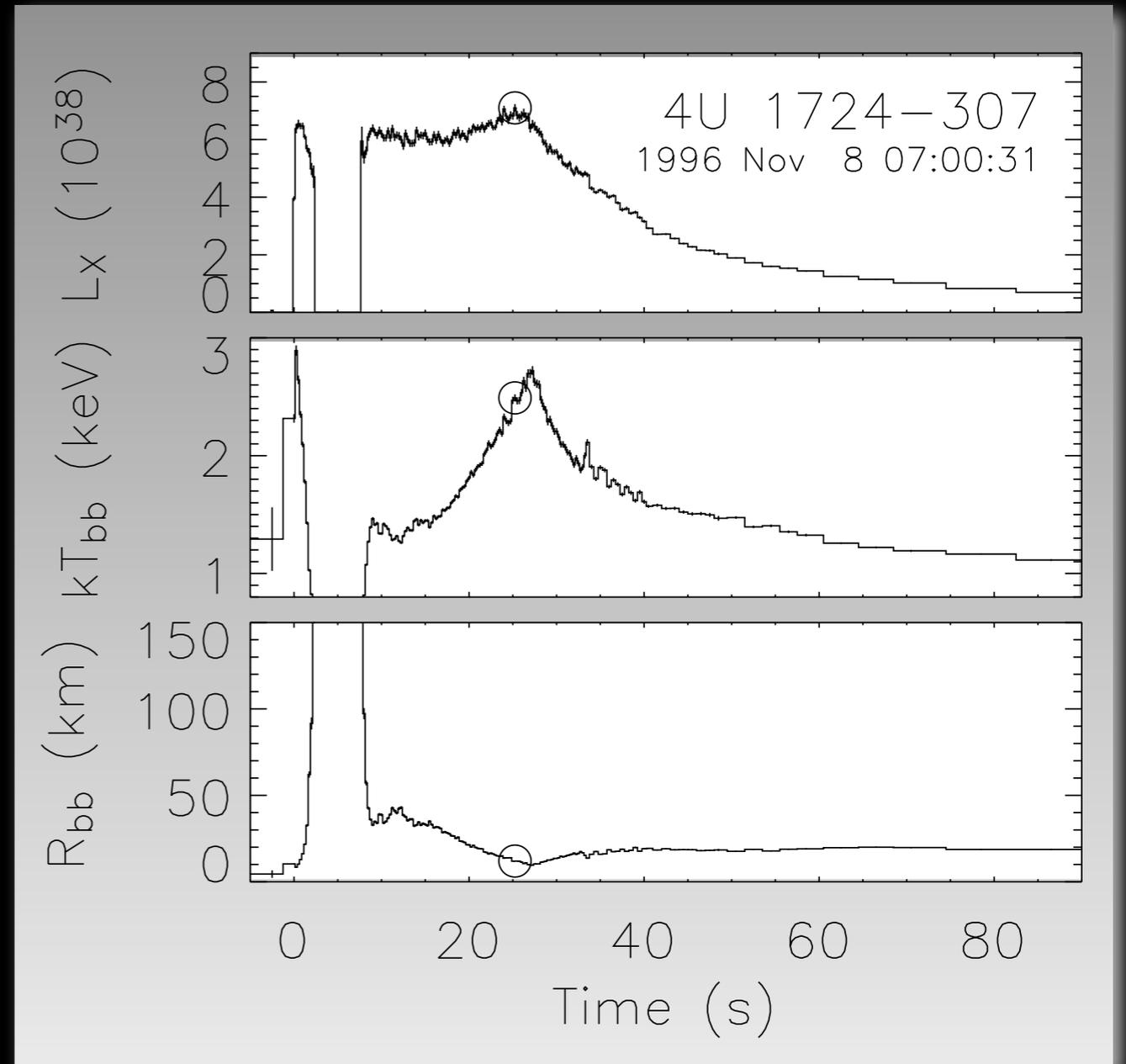
A sample of 1187 X-ray bursts from 48 sources



Eddington Limit



RXTE observations; Galloway et al. '08

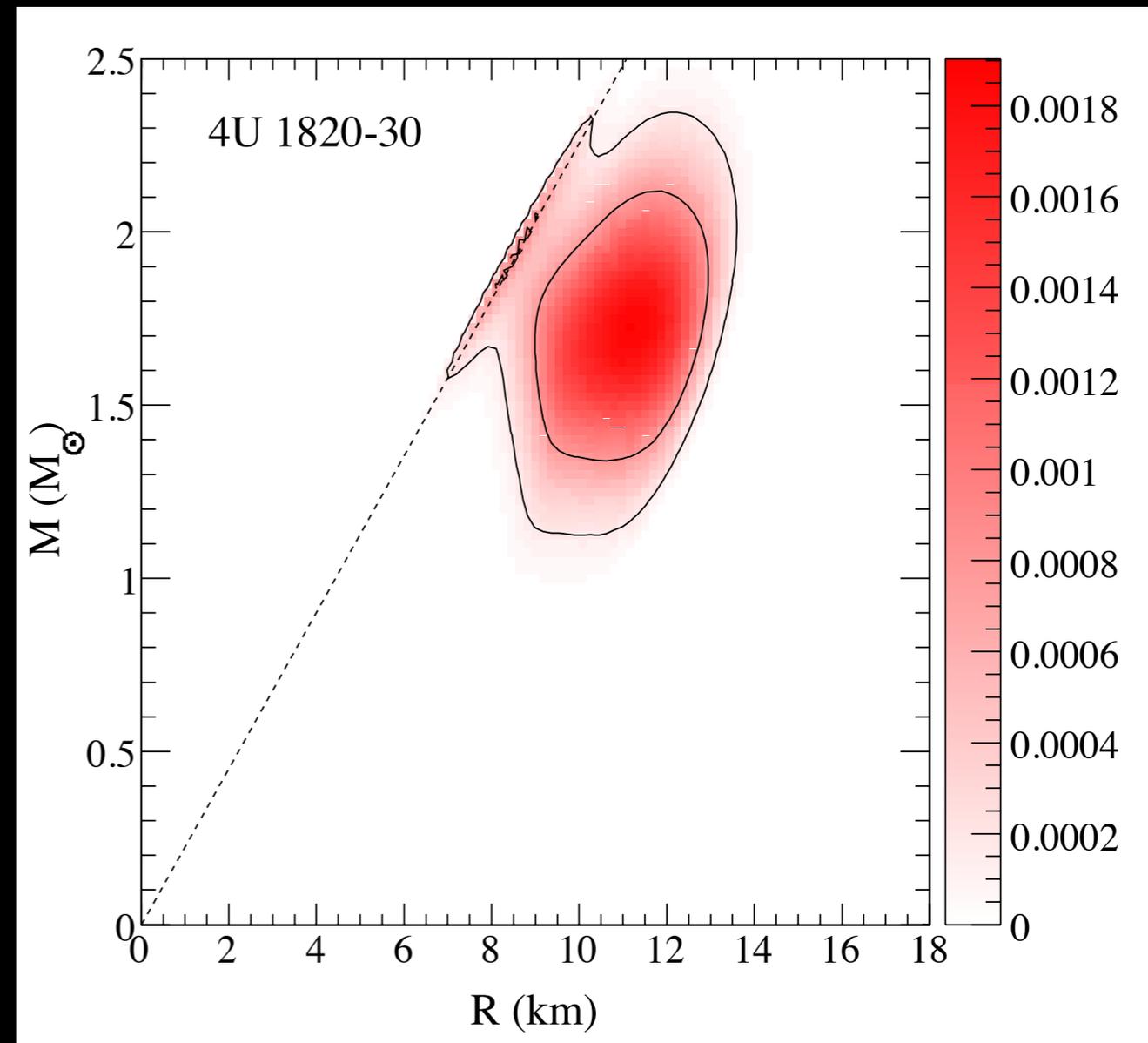


An Empirical Dense Matter Equation of State

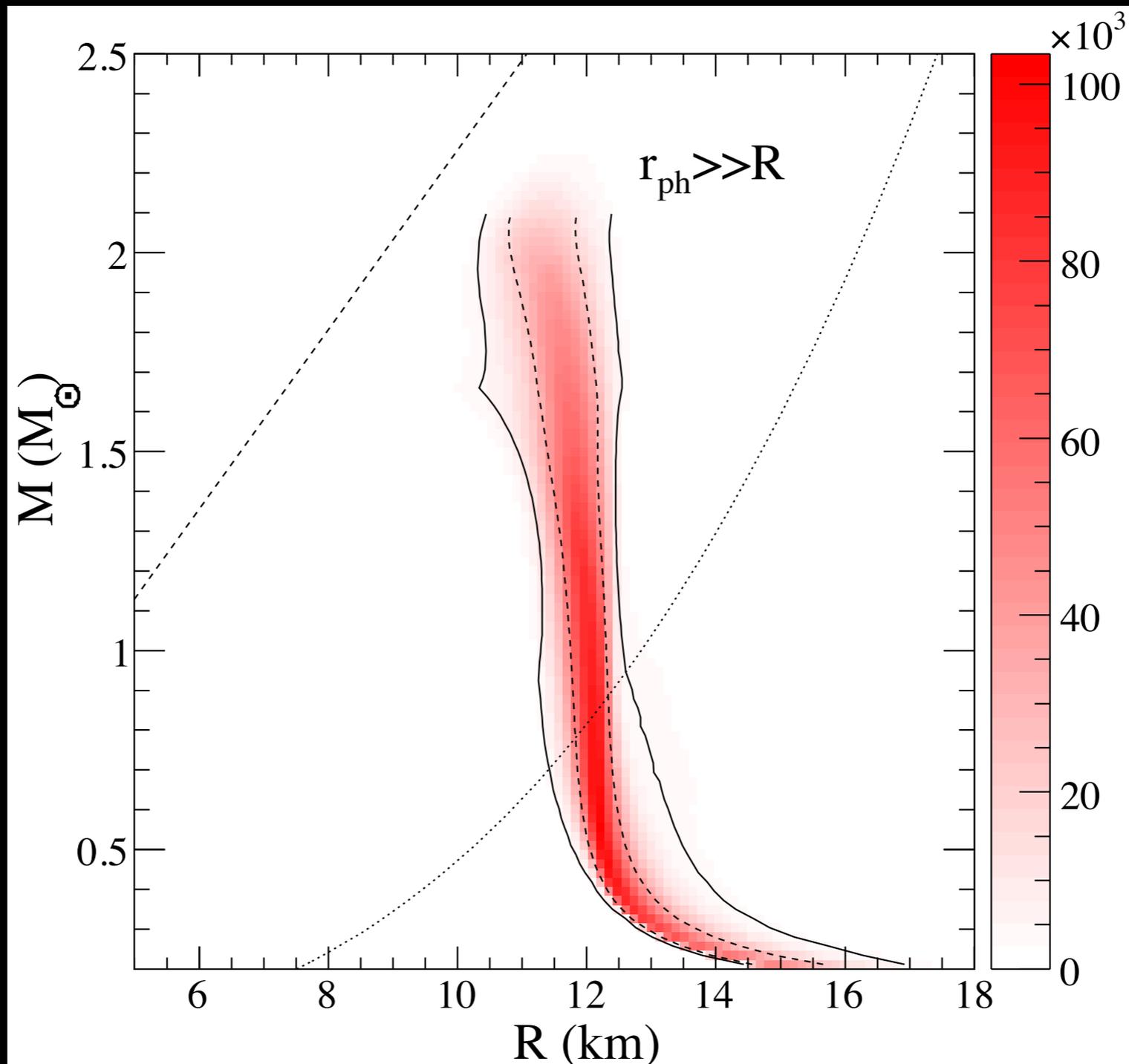
From X-ray bursts with *photospheric radius expansion* (van Paradijs, Özel et al., Steiner et al., Suleimanov et al.)

$$F_{\text{Edd}} = \frac{GMc}{\kappa D^2} \left(1 - 2\frac{GM}{r_{\text{ph}}c^2}\right)^{1/2}$$
$$\frac{F}{\sigma T_{\text{eff}}^4} = \left(\frac{R}{D}\right)^2 \left(1 - 2\frac{GM}{Rc^2}\right)^{-1}$$

Talk by A. Steiner

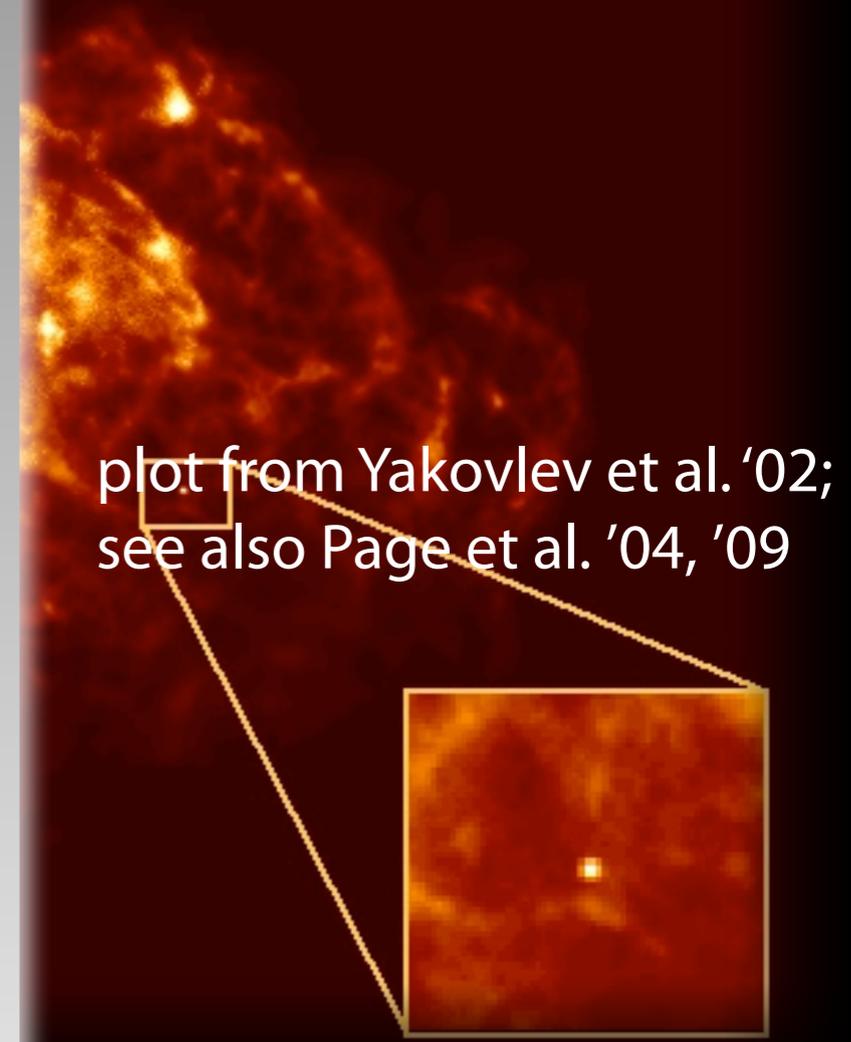
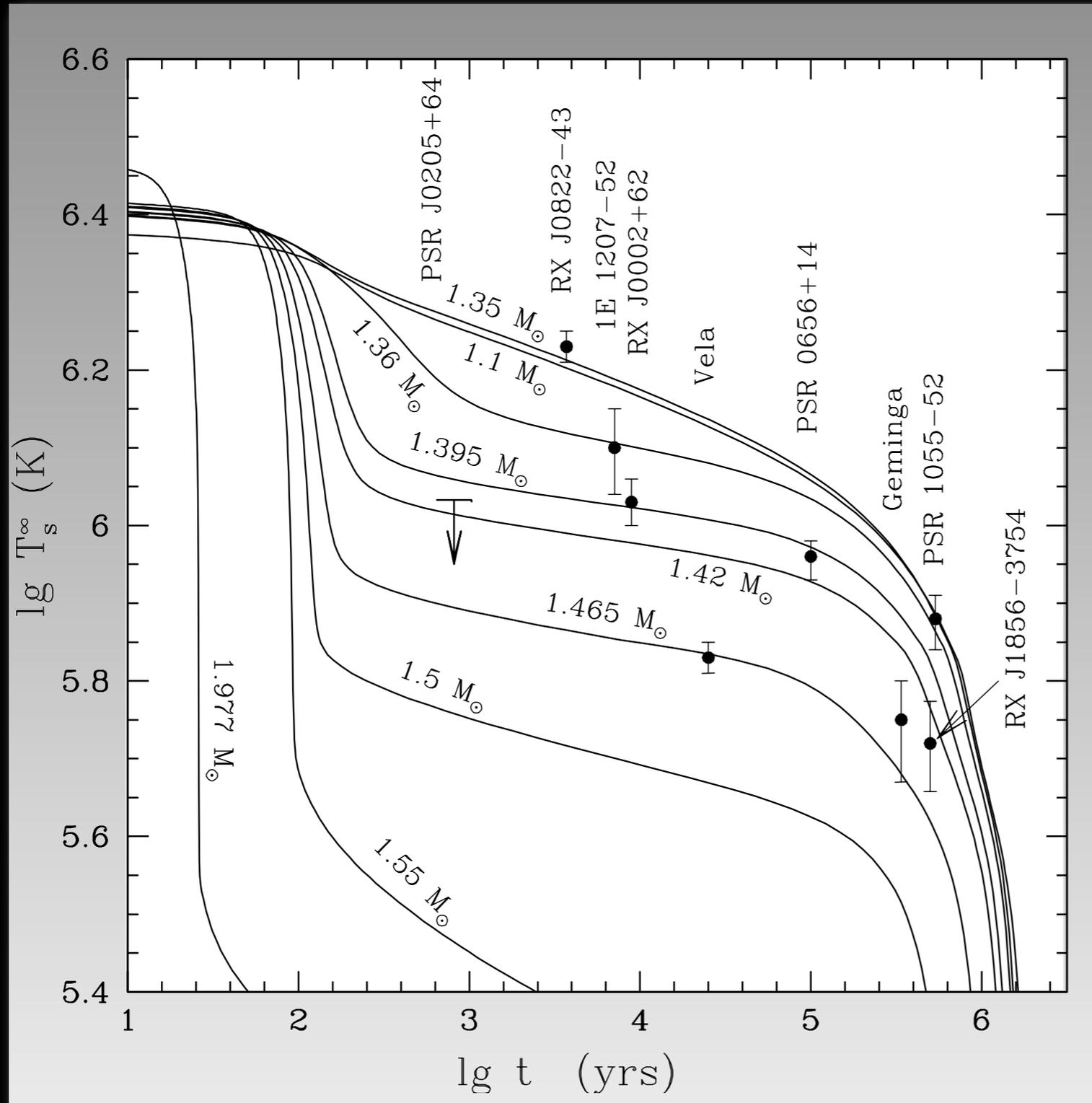


Steiner et al.; data from Guver et al. '10



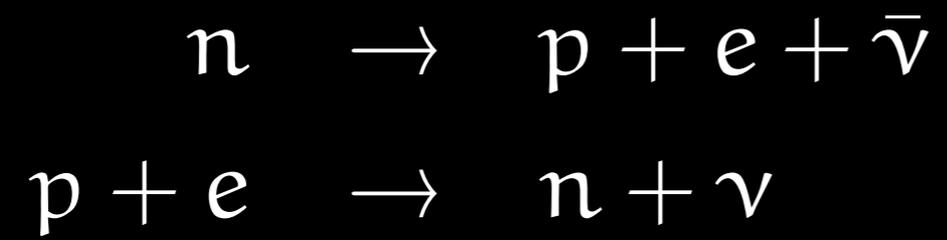
An empirical M - R relation (Steiner et al.)

Neutron stars cooling from their fiery birth



cooling: the Urca process

Gamow & Schoenberg 1941



but this is blocked...

Chiu & Salpeter, Bahcall & Wolff

mom. cons.

β -eq.

charge neutrality

$$p_{Fn} < p_{Fe} + p_{Fp},$$

$$\mu_e = \mu_n - \mu_p,$$

$$n_e = n_p$$

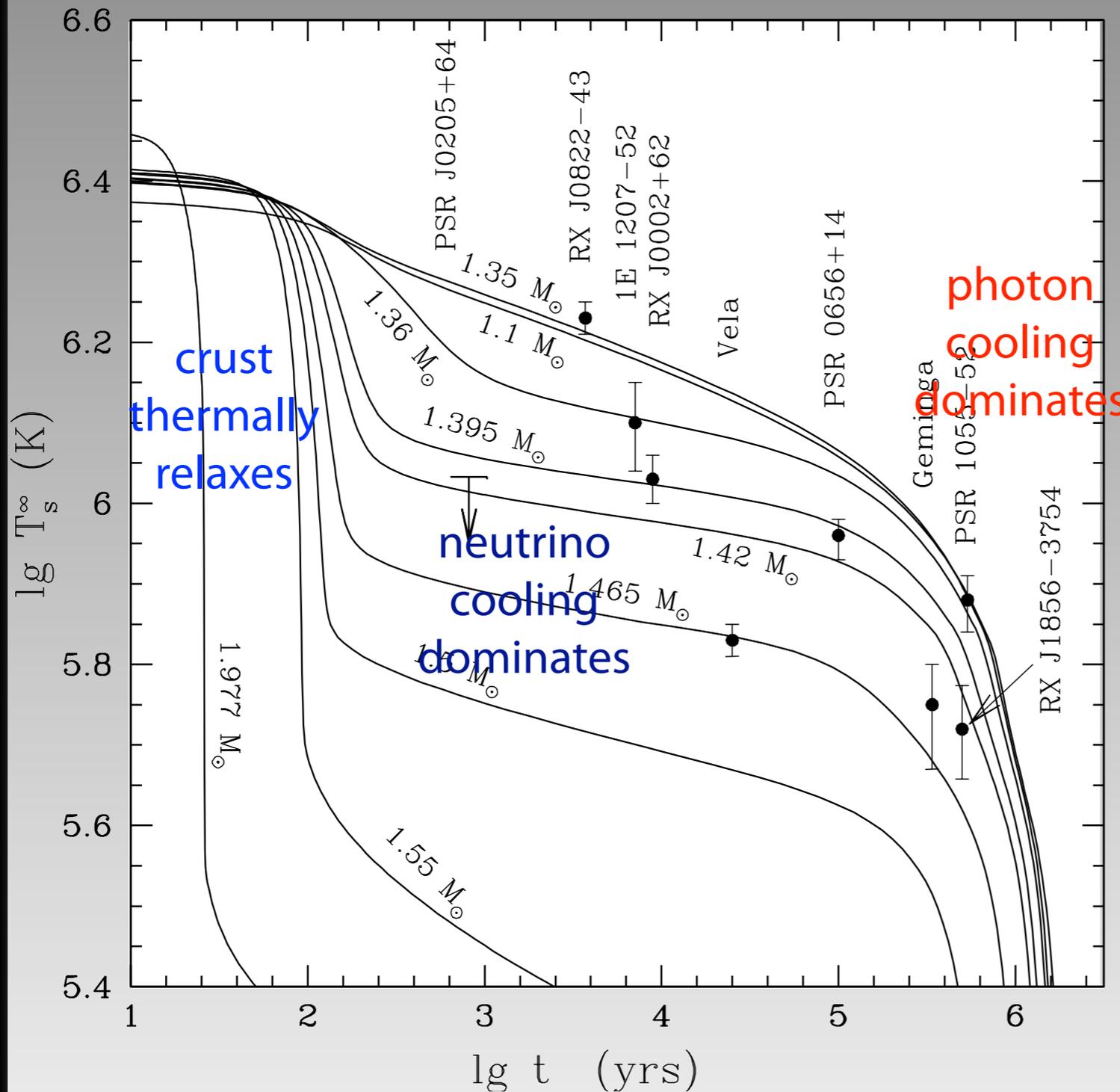
...unless $n_p/n > 0.11$

a high symmetry energy
implies that neutron stars
should cool rapidly!

—Lattimer & Prakash 2007



Neutron stars cooling from their fiery birth

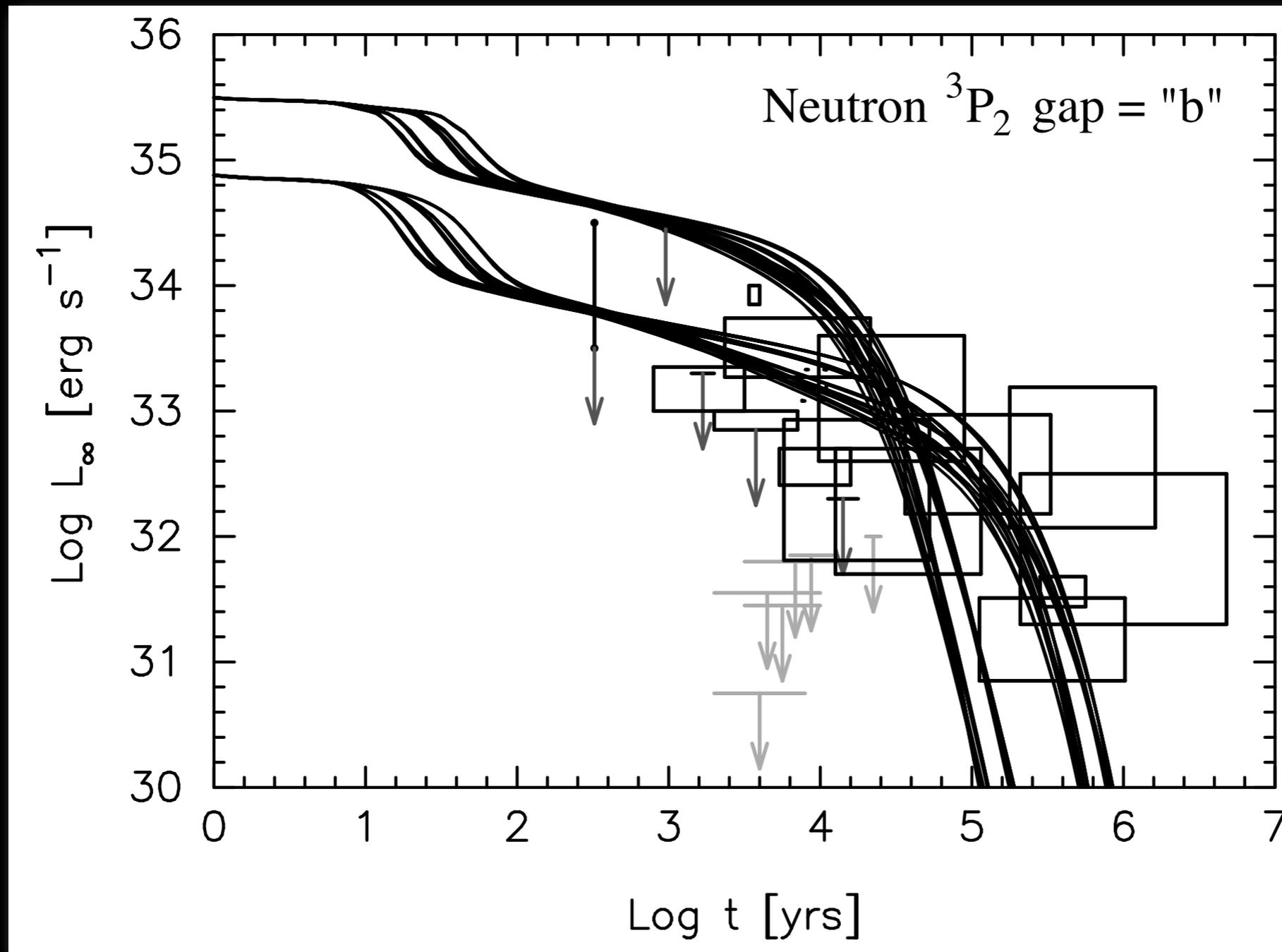


plot from Yakovlev et al. '02;
see also Page et al. '04, '09

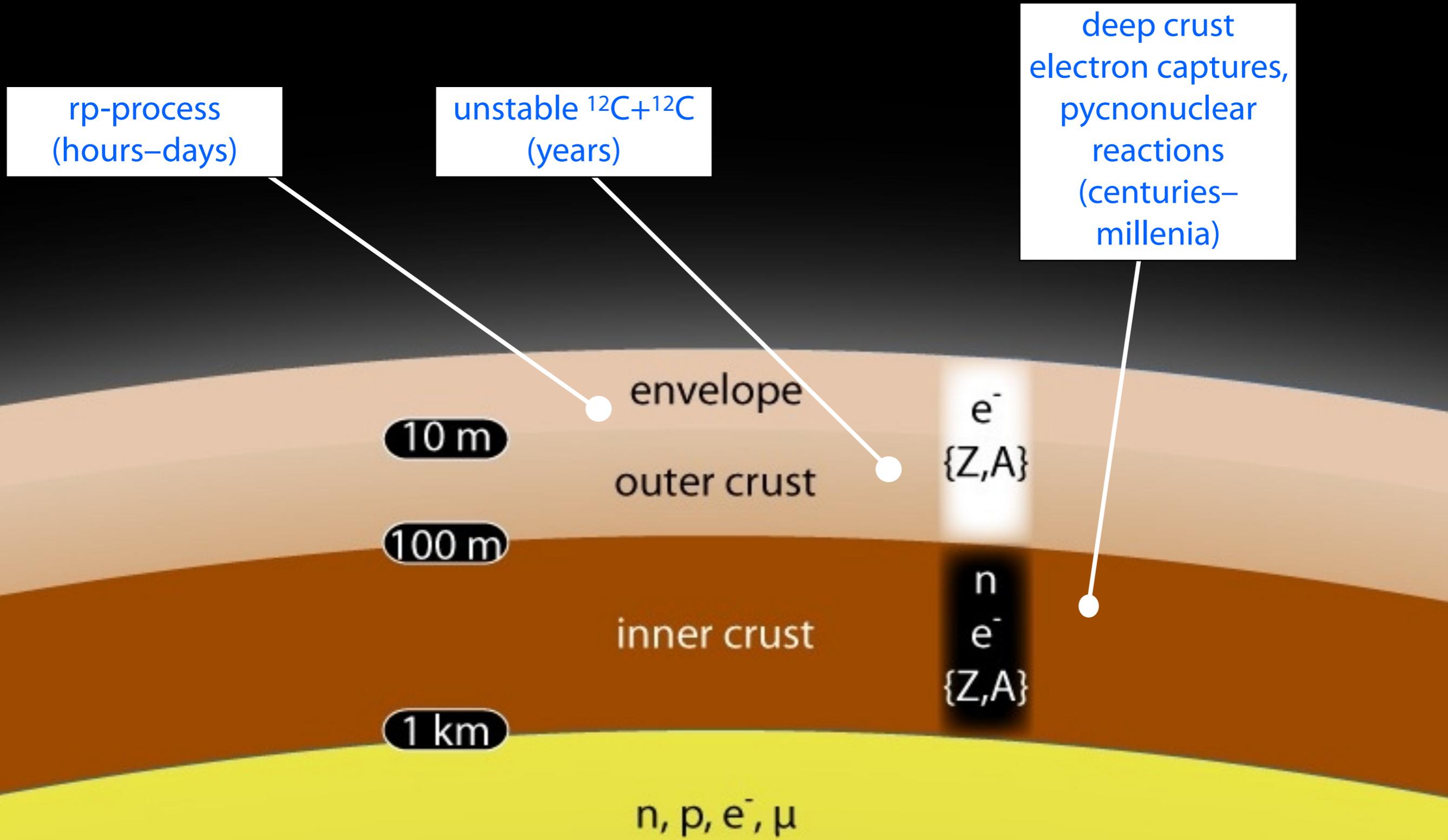
enhanced cooling (dUrca)
leads to rapid drop of
surface temperature after
crust has relaxed

Isolated neutron stars: no enhanced cooling in models

Page, Lattimer, Prakash, & Steiner 2009



accreting neutron stars: deep crustal heating

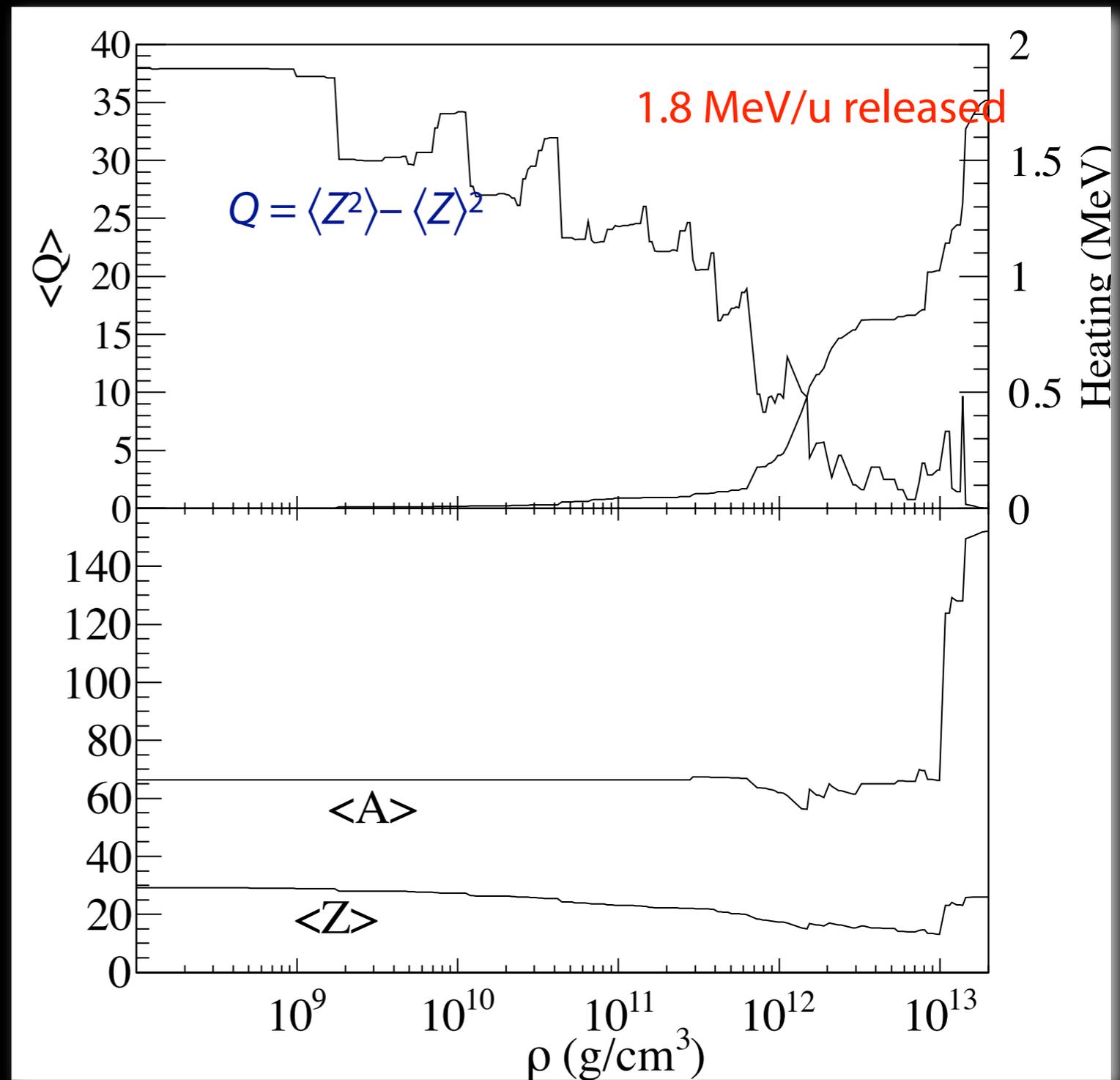


deep crustal heating

crust reactions deposit ~ 1.8 MeV/u
in the inner crust

1. core temperature set by
balance of heating, neutrino
cooling
2. crust is not in thermal
equilibrium with core

plot courtesy A. Steiner

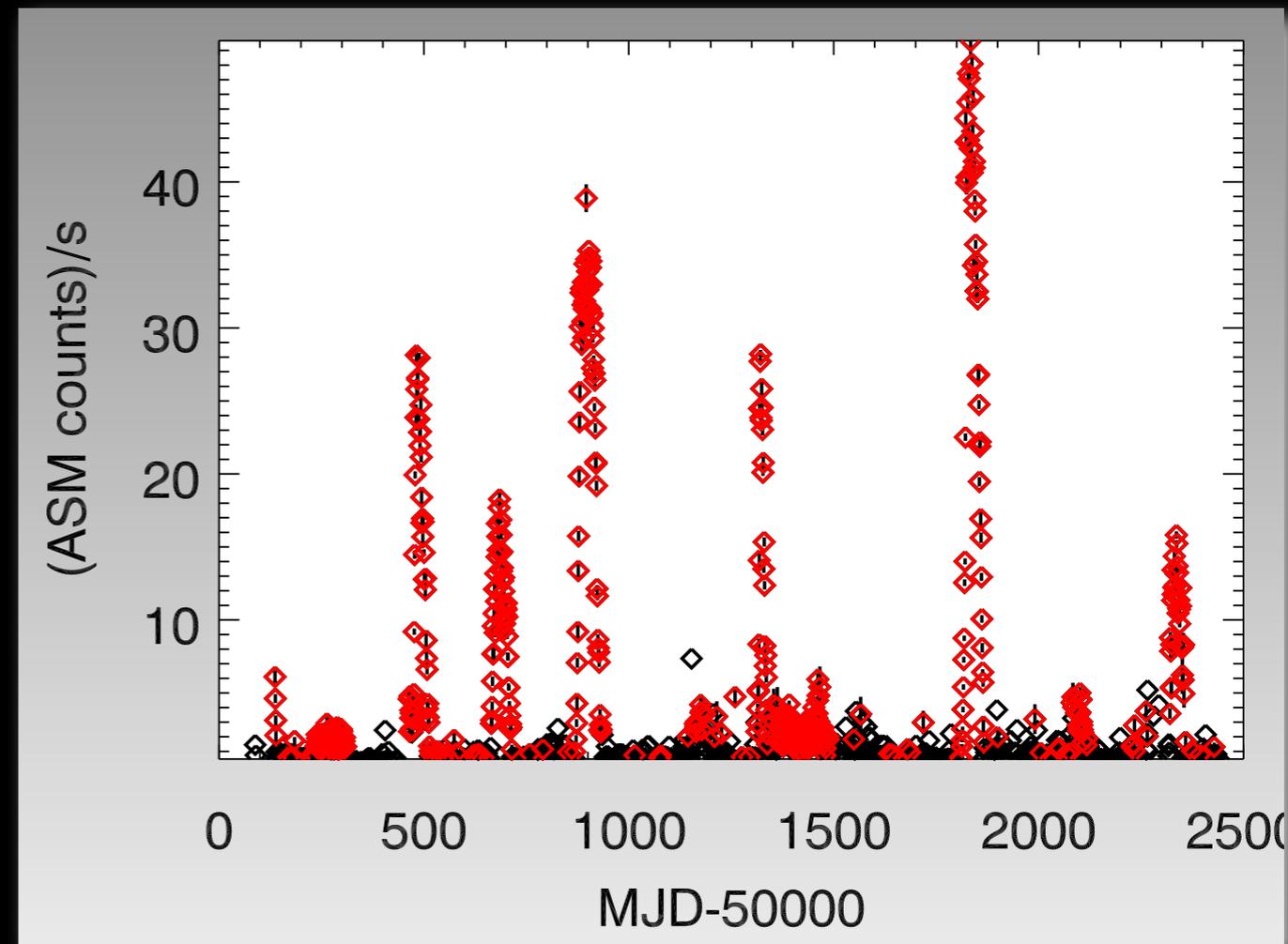


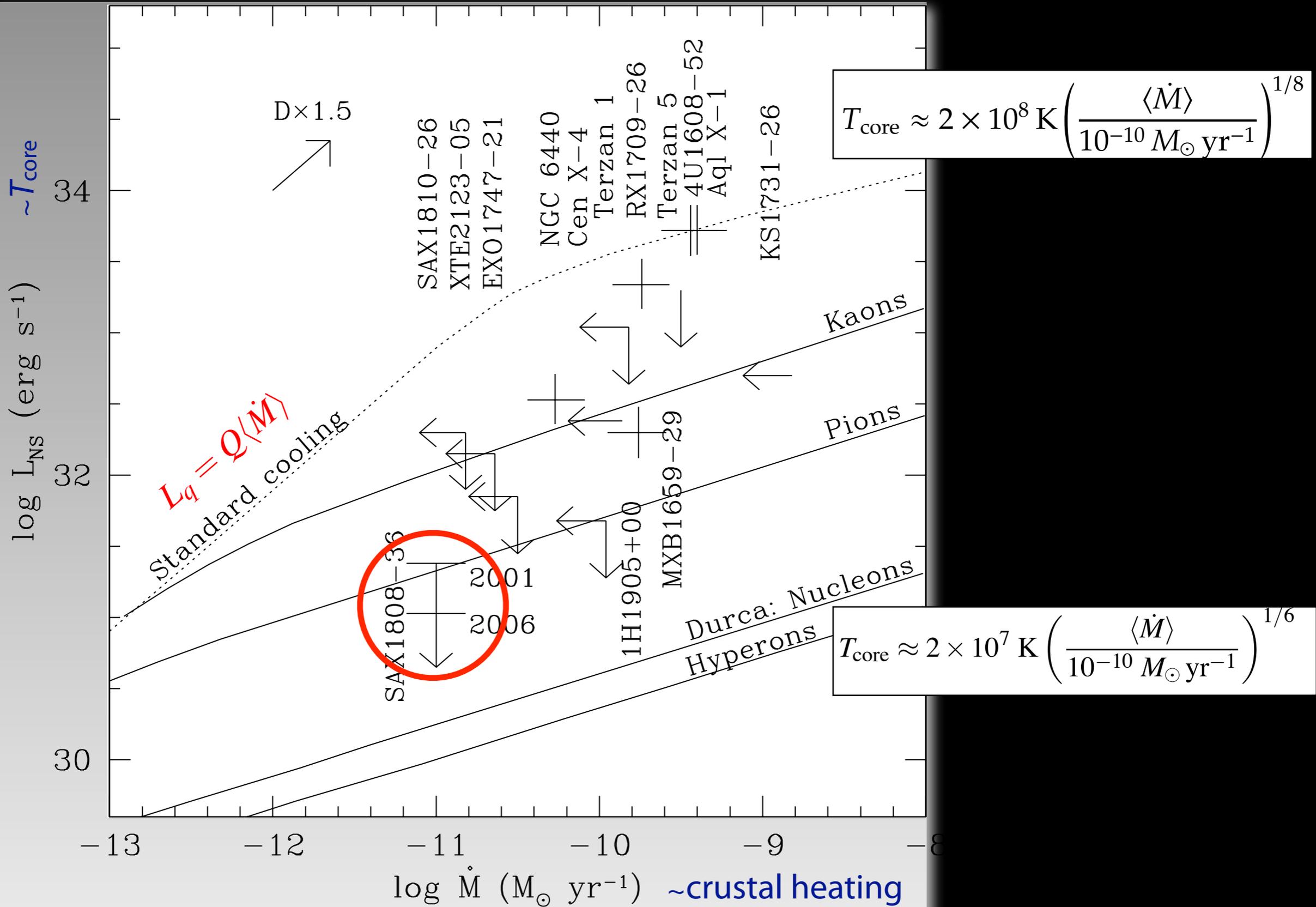
transients

Brown, Bildsten, & Rutledge; Colpi et al.; Yakovlev et al.

Many neutron stars accrete
intermittently
crust heats during outburst
crust cools during quiescence

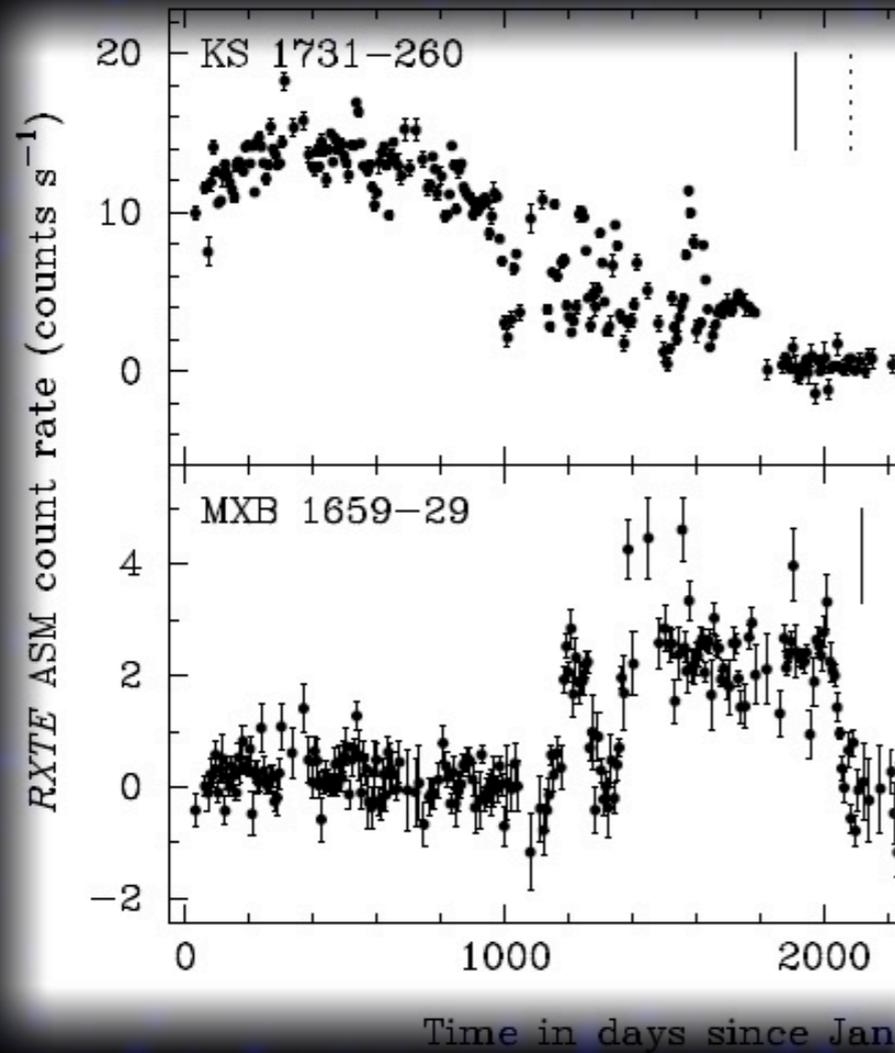
Aql X-1 lightcurve. *RXTE/ASM*



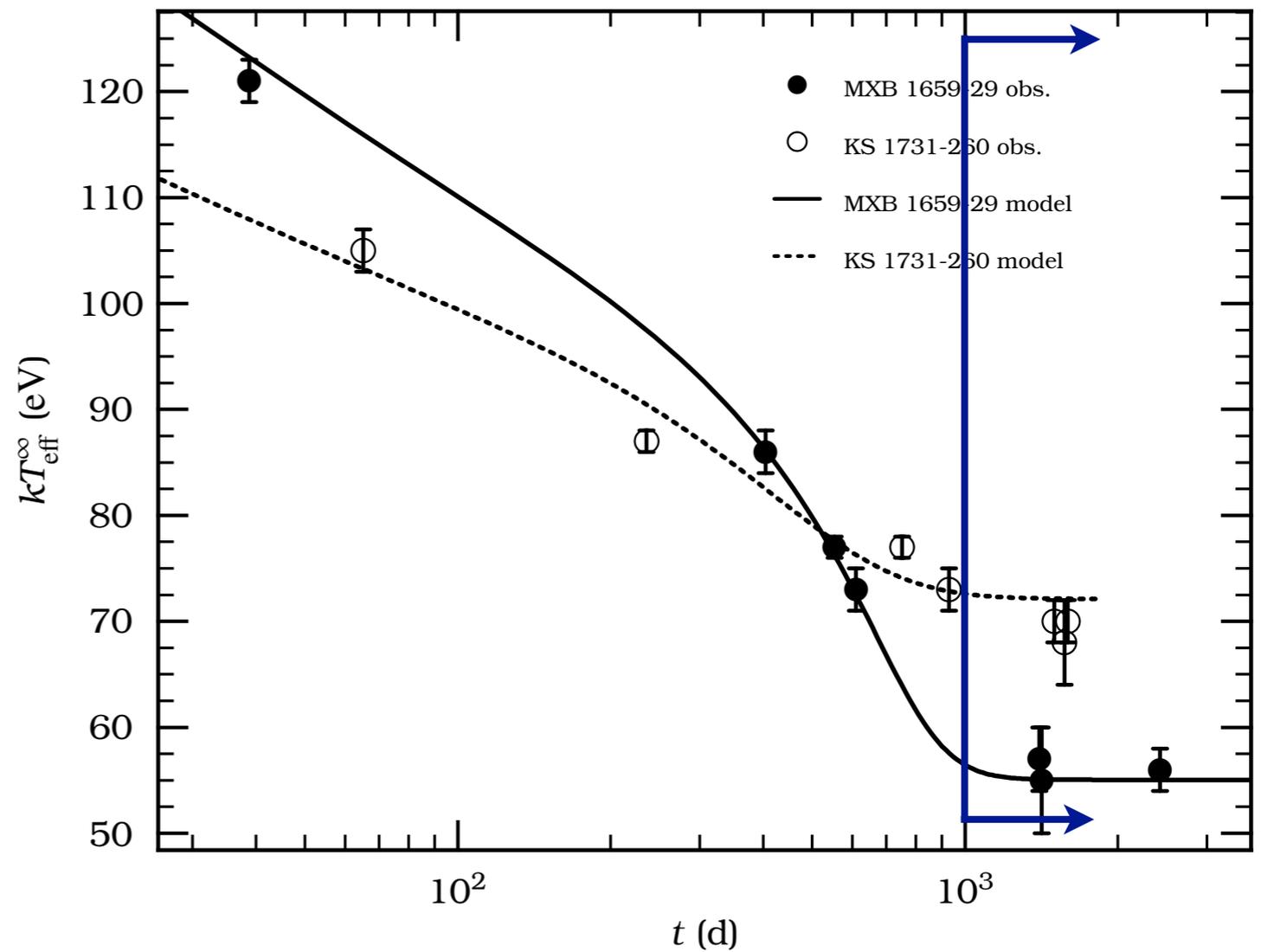


quasi-persistent transients

Rutledge et al. 2002, Shternin et al. 2007, Brown & Cumming 2009



data from Cackett et al. 2008
fits from Brown & Cumming 2009



Summary

A number of observational probes of the nuclear EOS are available

three examples

pulsar masses

masses and radii from X-ray bursts

cooling of isolated neutron stars, thermal relaxation of accreting transients

this is not inclusive: there are others

No single observation is ideal—it's astrophysics; but

These observations, taken together, offer interesting constraints and complement theoretical and experimental efforts

The best is yet to come!